

ANALYSIS OF SPHERICAL TEXTURE ON CHARACTERISTICS OF HYDRODYNAMIC JOURNAL BEARING

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Certificate

This is to certify that the thesis entitled, “**Analysis of spherical texture on characteristics of hydrodynamic journal bearing**” submitted by Mr. **Himalaya Dawani** to National Institute of Technology Rourkela, during the academic session 2013-2015 is a record of bonafide research work carried out by him under my supervision and is worthy of consideration for the award of the degree of Masters of Technology in Mechanical Engineering with specialization in **Machine Design and Analysis**. The embodiment of this thesis has not been submitted to any Other University and/or Institute for the award of any degree or diploma.

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Abstract

Texture produced on the surface of hydrodynamic journal bearing have certain changes in the coefficient of friction, load carrying capacity, wear rate, stiffness and damping of the two surfaces which are matting with each other. By creating texture or micro dimple on a hydrodynamic journal bearing, pressure increases and significant improvement has been achieved in its load carrying capacity, lubricant flow rate, coefficient of friction etc. In a present work numerical analysis has been carried out to determine the effect of using negative spherical surface texture on hydrodynamic journal bearing surface and it is compared with normal journal bearing i.e. journal bearing without texture. We can take any other profile also but spherical profile is easy to fabricate by laser technology or etching process. The Reynolds equation is solved numerically with the help of central finite difference method and analysis is done on the effect of texture height, asperity ratio and number of textures on the journal bearing. Here we also determine the behavior of parameters of hydrodynamic journal bearing (texture and without texture) with increase in eccentricity ratio. Texture journal bearing has a great importance in journal bearing because in vertical journal bearing initially it provides converging part which helps in increasing load carrying capacity which is not present in plain journal bearing. The author believes that such detail analysis of texture on journal bearing surface to study different tribological behavior under different condition will help researchers around the world.

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Nomenclature

ρ : Density of lubricant (Kg/m³)

η or μ : Dynamic viscosity of lubricant (N-sec/m²)

h : Lubricant film thickness (mm)

\bar{h} : Non-dimensional film thickness = $\frac{h}{c}$

$$\bar{z} = \frac{z}{L}$$

C : Clearance (mm)

C_g : Texture height (mm)

\bar{C}_g : Non-dimensional texture height

R : Radius of bearing (mm)

r : Radius of journal (mm)

$\bar{\theta}$: Non-dimensional co-ordinates

\bar{W} : Non-dimensional load = $\frac{W}{\frac{u\mu R^2 L}{c^2}}$

\bar{q} : Non-dimensional flow = $\frac{q}{c^3 \omega}$

\bar{F} : Non-dimensional frictional force = $\frac{F}{RL \frac{u\mu}{c}}$

\bar{f} : Non-dimensional co-efficient of friction = \bar{F} / \bar{W}

\bar{p} : Non-dimensional pressure = $\frac{pc^2}{6\eta\omega R^2}$

THR: Texture height ratio

AR: Aspect ratio

\bar{h}_b = Non dimensional texture height

k: stiffness (N/m)

\bar{k} : Non dimensional stiffness = $\frac{kc^2}{\mu\omega R^3 L}$

b: damping (N-m/sec)

\bar{b} : Non dimensional damping = $\frac{bc^3}{\mu R^3 L}$

\bar{p}_0 = Non dimensional perturbed pressure = $\frac{p_0 c^2}{\eta\omega R^2}$

\bar{p}_x = Non dimensional perturbed pressure = $\frac{p_x c^3}{\eta\omega R^2}$

\bar{p}_y = Non dimensional perturbed pressure = $\frac{p_y c^3}{\eta\omega R^2}$

\bar{p}_x^\bullet = Non dimensional perturbed pressure = $\frac{p_x^\bullet c^3}{\eta R^2}$

\bar{p}_y^\bullet = Non dimensional perturbed pressure = $\frac{p_y^\bullet c^3}{\eta R^2}$

Introduction

1.1 Background

Tribology is the branch of science which deals the surfaces, that are rub together or we can say that it is a scientific and systematic method to deal with interacting surfaces so that characteristics of the system can be improved. Tribology contains the study of chemistry, physics and mechanics of rubbing surfaces which include friction wear and lubrication of materials. From the ancient civilization various attempts have been made to minimize friction and wear so that we can transport men and materials from one place to another place in an economical way. This was the biggest challenge faced from past till now. In the nineteenth and the twentieth century lubrication came into picture which is usually put between interacting surfaces to minimize friction and wear.

Tower (1883) [1] who was a railroad engineer at that time has done a series of experiments on lubrication so that he could minimize friction to control the wear, for this he drilled a hole in a bearing so that lubricant (oil) can be poured through it, when shaft starts rotating he observes that the oil is coming out of that hole, to prevent this he put a plug on that hole. However, he noticed that plug is been thrown out from that hole, from there he concluded that oil film is being formed in between shaft and bearing that generates sufficiently high pressure to through the plug out of the hole. At the same time Petroff (1883) [2] was interested in calculating friction in journal bearings. For this he conducted some experiments and came out with some relationship between frictional force and operating parameters of bearing. However, he didn't notice that the oil film also generates the pressure which was given by tower hence he should be given the credit for

enhancing the concept of hydrodynamic fluid film lubrication. With the help of these developments, Reynolds (1886) [3] formulated the concept of formation of hydrodynamic fluid film lubrication in journal bearings and developed Reynolds equation for hydrodynamic lubrication. Reynolds equation also explained that hydrodynamic pressure which is generated between shaft and bearing is due to converging wedge shaped film, viscosity of oil and surface motion. It was also suggested by him that both pressure gradient and pressure for divergent portion will be zero at the film rupture boundary. The results gave by Reynolds were approximate as he didn't go for the integration of Reynolds equation. Sommerfield (1904) [4] obtained the analytical expression for pressure distribution, load carrying capacity, frictional force etc. by integrating the Reynolds equation. These equations form a basis later to design different types of bearings which was later used in different machineries. After these developments Kingsbury (1913) [5] developed gas bearing and showed that air film can also be used to carry load.

Current research target to modify the surface of journal by introducing texture on it and a detail study was done to find non dimensional pressure profile, non-dimensional load carrying capacity, coefficient of friction parameter, stiffness and damping coefficients by varying eccentricity ratio, asperity ratio and texture height.

1.2 Thesis overview

Creating surface texture or micro asperity or undulation on the surface the surface of mechanical component produces more wedging action as compared to the normal i.e. without textured bearing which will results in increase in pressure due to which there will be increase in load

carrying capacity, it also reduces friction because each undulation will act as a reservoir, stores lubricant and will provide it during starving condition. Micro asperities or texture can be of different shapes such as cylindrical, elliptical, spherical, circular, semi spherical and various others but we have done work using spherical micro asperity and it can be of different sizes also, so an optimum size should be used. Here we have done a theoretical analysis on hydrodynamic journal bearing by applying texture on it and find its effect on tribological parameters of hydrodynamic journal bearing.

1.3 Objectives of present work

In present research work we study how negative spherical micro asperities affect the properties of journal bearing. Following are the main objectives of present work:

- a.** Development of numerical solution to determine the effect of negative spherical micro asperity on hydrodynamic journal bearing.
- b.** Analysis is done by varying epsilon, asperity ratio, and texture height.
- c.** Comparison of tribological parameters for plain and textured journal bearing on varying eccentricity ratio, asperity ratio and texture height.

1.4 Basic concepts

1.4.1 Surface topography

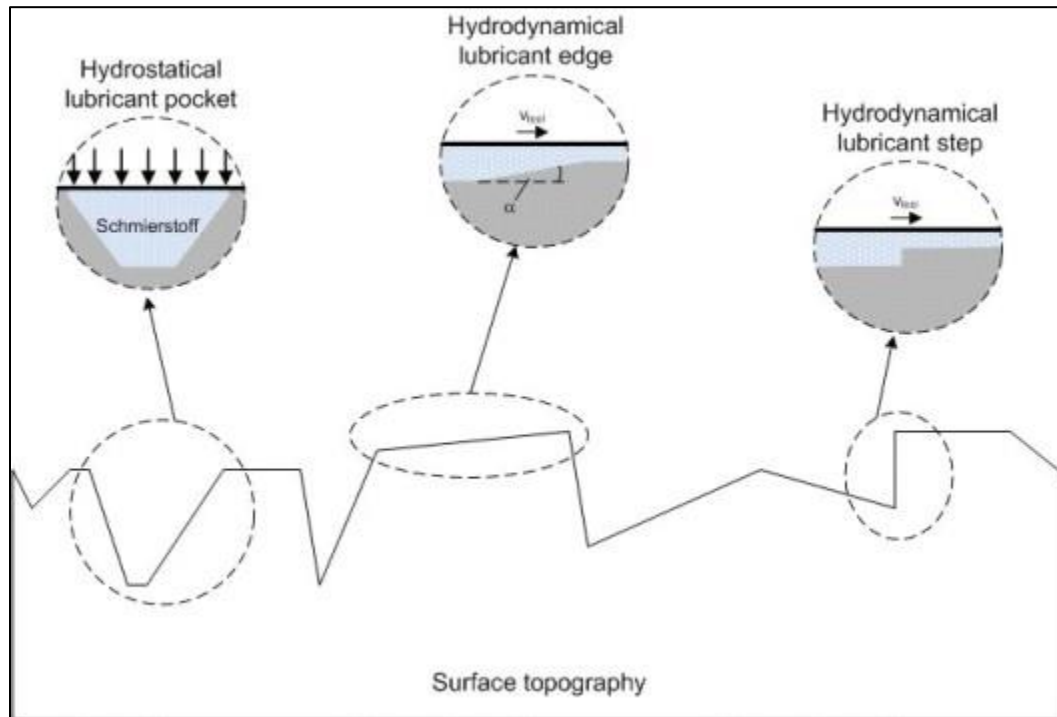


Fig 1.1: Surface topography (www.ptu.tu-darmstadt.de)

Usually due to mating surfaces tribological phenomenon occurs, so the variation in the surface texture is the key point for studying the mechanism of friction and wear. The imperfections at the surface at a nano level are measured and compared by macroscopic deviations from flatness. In general we consider all the surfaces to be rough except cleaved faces of mica. Roughness means peaks and valleys on the most part of the surfaces. The rough surface profile is always random unless and until some deliberately rough structures are made on it.

Most of the real engineering surfaces generally consist of random and non random features. The prime example of non random topological characteristics is a series of grooves formed by a shaper on a metal surface. The shaped surfaces also give rough texture because it also contains

high degree of random surface features.

These micro asperity or undulations have various shape and sizes. Their heights vary from fraction of a nanometer to several millimeters. Depending on the scale, they can be called macro deviations, roughness or waviness. Macro deviations are mainly recognized by asperities of small height with gentle slopes similarly for waviness the ratio of distance between asperities to asperity height is more than 40.

1.4.2 Friction

In many types of machinery rough surfaces generally rub against each other under an unlubricated environment which results in high friction and lubrication. Leonardo da Vinci (1452-1519) [6], Coulomb (1785) [7], Amonton (1699) [8] enunciates the fundamental laws of friction which are

- i. Frictional force varies with the normal load.
- ii. The force of friction is independent of the contact area.
- iii. Static friction is much larger as compared to kinetic friction.
- iv. Friction doesn't depend on the sliding velocity.

When the shaft rotates inside the bearing generally interlocking of asperities takes place due to which resistance to relative motion occurs which is the main cause of friction between two surfaces. Molecular attraction as the main reason for friction between two interacting surfaces is proposed by Hardy (1936) [9]. Bowen and Tabor (1950) [10] also proposed the theory in which the theory of cold welding, shearing, ploughing of the soft metal by hard asperities are the main reasons of friction. Tomlinson (1929) also gave the molecular attraction theory in which he

attributed frictional resistance to rupture of molecular bond between the two surfaces which are interacting with each other. For metal to metal contacts Bowen and Tabor theory is widely accepted and for polymer molecular bond rupture theory is widely used.

1.4.3 Wear

Wear can be defined as a removal of material from one of the surface of a body (which are interacting with each other) occurring as a result of relative motion at the surface. Generally tribological pairs are supplied with lubricant so as to avoid the excessive wear and damage which would be present if two surfaces interacted with each other in dry state, to reduce their frictional resistance to motion. The economic consequences of wear are many more such as:-

- Cost of replacement of parts
- Cost involved in machine downtime
- Loss of production.

The wear rate of rolling or sliding contact is often defined as the volume which is lost due to wearing of surface per sliding distances. For particular dry or unlubricated sliding situation the wear rate depends on various factors such as –:

- I. Normal load,
- II. Relative sliding speed,
- III. Initial temperature
- IV. Thermal, mechanical and chemical properties of the materials in contact.

1.4.4 LUBRICATION

There are various lubrication system types which are classified as [11]-:

Class I — Bearings in which lubricant is provided from an external source (e.g., oil, grease, etc.).

Class II — Bearings which have lubricant inside the walls of the bearing (e.g., graphite, bronze, etc.). Generally these bearings required lubricant from outside for achieving maximum performance.

Class III — Bearings in which materials of bearing is similar to that of lubricant. These are also known as “self-lubricating” and don’t require external lubricant.

Plastic bearings and oilites made from polyacetal are the examples of second type of bearing; metalized graphite bearings and PTFE bearings [11] are the examples of third type of bearings.

Plain bearings have plain inner surface, some of them are grooved which helps the lubrication in entering the bearing to cover the whole journal [12].

Bearings which have lubricant enclosed within the bearing walls are known as self-lubricated bearings. These are of many types. Bearings having porous walls are the first type of self-lubricated bearings which is known as sintered metal bearings. These porous walls use capillary action to draw oil [13] and discharges when heat or pressure is applied [14], for e.g. in self-lubricating chains. Second one is a solid one-piece metal bushing having a figure eight groove channel on the inner diameter which is filled with graphite, eight groove channel is there so that bearing can be lubricated inside and out [15] and its final form is a plastic bearing, in which lubricant is shaped into the bearing. The lubricant is discharged as the bearing is run in [16].

Types of lubrication can be classified as:

- I. Boundary condition.
- II. Full film condition.
- III. Dry condition.

1.5 Fluid lubrication

Fluid lubrication results in a boundary or full film lubricant mode condition. It can be lubricated in two ways i.e. hydro statically or hydro dynamically lubricated. Hydrostatically lubricated bearings are those bearing which use external pump to lubricate the bearing and constantly generate static pressure whereas the bearing which keeps the pressure in the oil film by the rotation of the journal is known as hydro dynamically lubricated journal bearing. State of hydrostatic bearings changes in to hydrodynamic state during rotation of journal [17]. Mostly oil is used by hydrostatic bearing whereas hydrodynamic bearing can use grease or oil.

"Oil whirl" is the disadvantage in high speed machinery in fluid lubricated bearing. The main reason of oil whirl occurrence is lack of stability in lubrication wedge, small disturbance in the journal cause reaction forces from the oil film which results in further movement, causing both the oil film and journal to "whirl". Generally the frequency of whirl is about 42% of the turning speed of journal. In worst cases oil whirl may cause the direct contact between the journal and the bearing, due to which bearing wears out rapidly. In some of the cases the frequency of whirl matches with the natural frequency this leads to catastrophic failure [17, 18] and it is very destructive.

This can be prevented by a stabilizing force which is applied to the journal. There are number of bearing designs which used bearing geometry to either provide an obstruction in the way of

whirling fluid or to provide a stabilizing load so as whirl can be minimize. One such is called the *elliptical bore* or *lemon bore*. In this design, shims are installed between the two halves of the bearing housing and then the bore is machined to size. After the removal of shims, the bore resembles a lemon shape, which reduces the clearance in one direction of the bore and increases the pre-load in that direction. The main disadvantage of this design is its lower load carrying capacity, in comparison to typical journal bearings. It is also still capable to oil whirl at high speeds, however its cost is relatively low [17].

Another design is the *pressure dam or dammed groove* [19] which has a deep relief cut in the center of the bearing over the top half of the bearing. The groove suddenly stops in order to create a downward force to balance the journal. Design has a high load capacity and rectifies most oil whirl situations. The disadvantage is that it works only in one direction. Offsetting the bearing halves have the same effect as the pressure dam. The only difference is that the load capacity increases as the offset increases [17].

A more ultra-design is the *tilting pad* design, which utilizes multiple pads that are designed to move with loads which changes frequently. It is mostly used in very large applications such as modern turbo machinery because it almost completely cancels out oil whirl and saves the bearing from catastrophic failure.

1.6 DESIGN

The design of a bearing mainly depends on the type of motion which the bearing is provided. The three types of motions which are possible are:

Journal (friction, radial or rotary) bearing: It is a type of plain bearing; this type of bearing usually consist of a shaft rotating in a bearing [20] for e.g. in rail road car and locomotive applications a journal bearing specifically focus to the plain bearing once used at the ends of the axles of railroad wheel sets, enclosed by journal boxes [21].

Linear bearing: This type of bearing provides linear motion; it can take the form of a circular bearing and shaft or any other two similar surfaces (e.g., a slide plate) [20].

Thrust bearing: This bearing provides a bearing surface for forces acting axial to the shaft [3].

1.6.1 Integral

Integral plain bearings are developed into the object of use. It is a hole which has been prepared into a bearing surface. The materials used in this type of bearing are Babbitt or cast iron and a hardened steel shaft [21].

These bearings are not as common because in these types of bearings, bushings are easy to adjust and it can be replaced if necessary [20]. Looking with respect to material, an integral bearing may be lesser in cost but it cannot be replaced, if it wears out. If wearing occurs, then the item may be replaced or reworked for e.g. hinge, which is usually in both thrust bearing and a journal bearing.

1.6.2 Bushing

An independent plain bearing which is inserted into a housing so that it give a bearing surface for rotary applications is known as bushing, it is also known as a *bush*; and also the most common form of a plain bearing [22]. There are carious design of bushings which include sleeve, *split*, and *clenched* bushings. A sleeve, split, or clenched bushing is only a "sleeve" of material with an inner diameter (ID), outer diameter (OD), and length. These three bushing are different in a way

that a solid sleeved bushing is solid whole way around, a split bushing has a cut along its length, and a clenched bearing similar like the split bushing but with a clench across the cut. A flanged bushing known as sleeve bushing consists of a flange at one end which is extending radially outward from the outer diameter. Flange is used to position the bushing when it is installed or to provide a thrust bearing surface [23].

1.6.3 Two-piece

Two-piece plain bearings or full bearings are mainly used in industrial machinery [24], these are mainly used for larger diameters, for e.g. in crankshaft bearings. *Two-piece* plain bearings contain two half which are called shells [25]. To keep the shells positioned various systems are used. The method which is mostly used is a notch in the housing tab that matches with the parting line edge to restrain axial movement after installation. For shells which are large, thick button stop or dowel pin is used. The button stop is coupled to the housing, while the dowel pin keys the two shells together. Another method which is not generally uses a dowel pin that keys the shell to the housing through a slot or hole in the shell [26].

The distance between one parting edge and the other is larger than the corresponding distance in the housing so that pressure which is required should be small which keeps the bearing in place as the two halves of the housing are installed. Finally, the housing circumference is also merely smaller than the shell's circumference so that the bearing crushes lightly when the two halves are bolted together. Due to this large amount of radial force around the entire bearing will be developed, which keeps it spinning. It also forms a better interface for heat to travel in to the housing from bearing [25].

1.7 SURFACE TEXTURING

Surface texturing is a technique by which micro dimples or undulation or asperity is created on the surface of materials which improve tribological properties such as load carrying capacity, wear resistance, friction resistance etc. of the material [27]. These undulations or micro dimples act as a reservoir which feed the lubricant to contact surfaces during starved conditions reduces the friction and wear. If the surface is mostly rough then wear will develop quickly, however if most of the surface smooth, inadequate lubrication or seizure may occur.

Different types and shapes of grooves can be produced by various techniques which are-

- I. Lithography technique
- II. Etching of silicon wafer
- III. Mechanical indentation
- IV. Laser surface texturing
- V. Abrasive jet machining
- VI. Novel dressing technique

From above all techniques, laser surface texturing technique is generally used because of its cleanliness to environment and provides outstanding control on the properties of micro-dimple. It is known that friction and wear create losses but efficiency and lifespan can be extensively improved by using laser surface texturing.

Laser surface texturing mainly have three functions:-

- **Micro hydrodynamic bearing**:- During relative motion between two contact surfaces each micro dimple acts as a small hydrodynamic bearing. This hydrodynamic effect results from the pressure gradient which forms in each cavity. Shear forces acting on a

lubricant at the surfaces due to relative motion between two contact surfaces, which results in wedge formation and due to wedge formation pressure profile is developed. This type of bearing reduces friction and wear. For micro hydrodynamic bearing lubrication is necessary.

- **Lubricant reservoirs-:** Micro dimple act as a reservoir of lubricant which will feed lubricant to the contact surfaces whenever starved condition arises due to which friction and wear decreases which prolonged the life span of the material. Since these micro reservoirs exist as asperities or micro dimples on the surface of bearing, so the geometry and pattern of micro dimple must be closed so as to prevent the force flow of lubricant out of channel.
- **Debris trap-:** Micro dimple provide the space for debris to fit in to it and reduce the abrasive wear which was earlier taking place in contact zone.

1.7.1 Deterministic Asperities

Micro asperities or undulation are the small dimples on a surface of the metals that creates surface roughness. Deterministic micro asperities are the asperities of definite shape, size, and orientation. Macro asperities have specific surface features with very low height to diameter (aspect ratio) <0.0001 . Macro asperities are mostly few in numbers which are created on the surface of bearing with comparative ease by using various processes such as laser itching and grinding. Similarly deterministic micro asperities are smaller in average diameter, significantly large in number with larger aspect ratio. There are various methods to fabricate such asperities such as photo etching, laser texturing, LIGA process (German acronym for lithography, electroplating & molding) & UV photo lithography process. Modified LIGA process has UV light as a substitute for X-ray. Surface texturing concept was developed after 1990 using

primarily laser technology. This process found a large number of industrial applications like in case of mechanical seals this technology enhances axial stiffness. Surface texturing technique is an efficient technique in tribology field because it reduces friction coefficient, wear rate and increases load capacity without changing dynamic stiffness and damping coefficient.

Positive Asperities

These are surface features of micro type of any arbitrary geometry that is in the form of bumps on a surface. The protrusions (bumps, posts) are called positive asperities. Figure 1.2 and figure 1.3 shows the positive asperities.

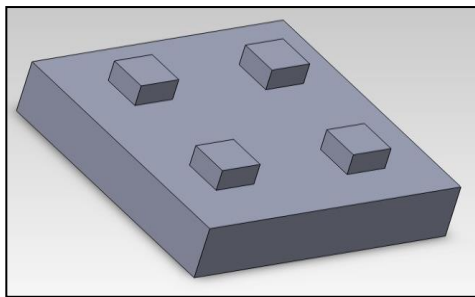


Fig 1.2 Square Positive Asperity

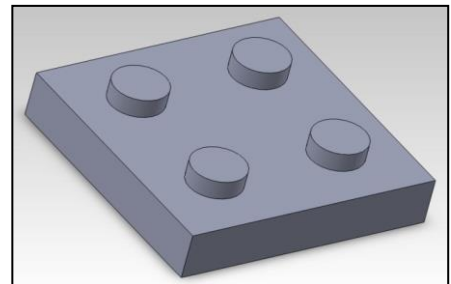


Fig 1.3 Cylindrical Positive Asperity

Negative asperity

These are the surface features of micro type of any arbitrary geometry that are in the form of cavities on a surface. These cavities (holes) are called negative asperities. Figure 1.4 and figure 1.5 shows the negative asperities.

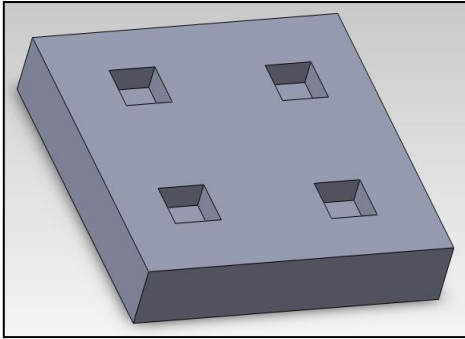


Fig 1.4 Square Negative Asperity

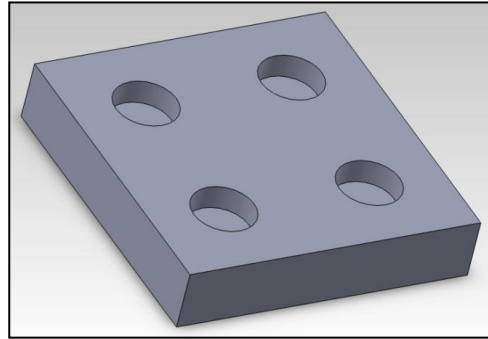


Fig 1.5 Cylindrical Negative Asperity

(a)-: We can produce both positive and negative dimples of various cross section and shape with the help of Lithography technique.

(b)-: Only way of producing spherical or conical geometry is by laser texturing method.

Chapter 2

Literature survey

Considerable amount work has been done on the surface texture on hydrodynamic journal bearing but still a lot of work has to be done further. This chapter contains a list of works which have already been done in the past on the texture profile journal bearing which helps in understanding various aspects of the surface texture journal bearing and study the performance parameter.

It is now easy to improve the tribological performance by surface texturing technique, which includes decrease in friction parameter, increases life, load capacity and lowers the energy consumed by producing the micro dimples on bearing surface. The technology of surface texturing is expected to be a very sufficient in future because it improves tribological properties of component as demonstrated by Priest [28].

Surface texture improves the tribological properties of boundary lubricated sliding surfaces. By producing micro dimples or asperities on the surface, we can improve the tribological properties of the component. After then lubricant is supplied inside the system with the help of small reservoirs, so as to reduce friction and increase lifetime of the system.

There are various situations of tribological contact in which improvement in characteristics has been seen by texturing such as in lubricated mechanical components e.g. aerodynamically lubricated magnetic hard disks, honed cylinder surfaces

Petterson and Jacobson [29] in 2003 determine the behavior of wear and friction properties of boundary lubricated textured surfaces. In this they determine how lubricant can be supplied to the interface and how to separate wear particles according to their sizes, shapes and orientation.

Series of surface textures of parallel grooves or square depression of various characteristics have been produced by various techniques such as anisotropic etching of silicon wafers and lithography.

Wakuda et al. [30] experimentally found that tribological characteristics of any system mainly relies on the size and density of micro dimples, and have very little effect of shape .

According to Siripuram et al. [31] there are various types of shapes of microstructures such as spherical, cylindrical, hemi spherical, triangular cross-section, square etc. and according to his experimental results square asperities gives the largest leakage rate and triangular asperities gives the smaller leakage rate.

Kovalchenko et al. [32] have studied laser texturing expanded the contact parameters in terms of speed and load of hydrodynamic lubrication. Laser surface texturing technique is more efficient at higher loads and speed provided that viscosity of lubricant should be high.

Mehenny et al. [33] had done a theoretical analysis to determine the influence of circumferential waviness of the journal in automotive engine on the lubrication of bearing. Assumption was made that surfaces of bearing and shafts have to be rigid and lubricant isoviscous.

Ronen et al. [34] provided the way to improve the tribological characteristics of the reciprocating automotive components by studying the micro-surface structure in the form of micro pores. He found that surface texturing on the surface of system can be used effectively to maintain hydrodynamic effects even with nominally parallel surface, and the optimum surface texturing improves the tribological characteristics in reciprocating automotive components.

Brizmer [35] demonstrated the use of laser surface texturing in parallel thrust bearing. It is found that the characteristics of textures may affect characteristics of bearing. The effects of texturing will be positive and it will improve the performances of journal bearing.

Values of friction torque, maximum pressure, minimum film thickness, axial film flow, and rupture film angle for both the surfaces (smooth and textured) are compared for different texture area fraction δ^2 . The following results are obtained:-

- With the increase in area of dimples, frictional torque will decrease for textured surface, whereas for smooth surface it will remain constant throughout.
- Maximum pressure will decrease with increase in surface area for textured surface, whereas for smooth area it will remain constant.
- Minimum film thickness increases with increase in area of dimples for textured surface, whereas for smooth surface it will remain constant.
- Axial film flow will decrease with area for textured surface whereas for smooth surface it will remain constant.
- Rupture film angle will increase with increase in area for textured surface, whereas it remains constant for smooth surface.

Friction torque, maximum pressure, minimum film thickness, axial film flow and rupture film angle are dissimilar from corresponding values which involve smooth surface. Difference becomes much larger as we increase texture area fraction.

Christensen [36], Tonder [37], Elrod [38], and Patir and Cheng [39] established the theories for hydrodynamic lubrication of rough surfaces. Surfaces of material can have different texture positions because of running-in and machining process. Christensen's concept [36] was extended

to the two sided roughness cases by Rhow and Elrod [40], Prakash [41], Zhang and Qiu [42, 43]. The averaged flow model of Patir and Cheng [44] was used by Boedo and Booker [45] for the analysis of isotropic Gaussian surface roughness effects in engine bearings. Christensen's stochastic model [36] was used by the Zhang and Qiu [42, 43] for determining the engine bearings effects of two-sided roughness with three types of roughness texture, in which the effect of transverse roughness only concerned the dimensionless bearing parameters.

The nominal minimum film thickness constantly decreases by the isotropic roughness. Minimum film thickness was negligibly affected by the transverse roughness. Increase in maximum film pressure is constantly seen due to transverse roughness, while decreases in nominal minimum film thickness is seen due to longitudinal roughness and both the roughness's increase or decrease the maximum film pressure is closely related to characteristics of nominal geometry and operating factors.

Asymmetric pressure distribution in the dimple area provides extra hydrodynamic pressure to separate the surface, due to inertia of lubricant. According to Tonder [37], if series of micro dimples are produced at the starting of sliding surface, it will generate extra pressure and will support higher load. Brizmer et al. [46] concluded that partial texturing generates collective dimple effect which results in pressure distribution over textured zone. Etsion and coworkers [47] concluded that dimple depth is most important factor for load generation. According to Yu and Sadeghi [48] there is some optimum value at which bearing can take maximum load. Sachlin et al. [49] states that maximum load is achieved when recirculation occurs at certain groove depth. According to Arghir et al. [63] if we analyze a single cell having micro dimple showed that, anti-symmetrical profile is induced due to inertia effects which results in positive load on the upper sliding surface.

Theoretical and numerical analysis

Generalized fluid equation & Reynolds's equation-:

The equations used for the generalized journal bearing is derived from the simplification to the general fluid mechanic equations governing the conservation of mass (continuity equation), momentum (Navier-Stokes equations), and energy (energy equation). The theory of hydrodynamic bearing on a differential equation was derived by **Osborne Reynold**. From Reynolds' theory we get the mechanism of lubrication through the generation of a viscous liquid film between two interacting surfaces. There are two conditions for the occurrence of hydrodynamic lubrication.

- There must be relative motion between two interacting surfaces with velocity which is sufficient for a load carrying lubricating film to be generated.
- Surfaces which are interacting must be inclined at some angle to each other, i.e. if they are parallel, pressure field will not be formed in the lubricating film and required load can't be supported.

Reynolds's equation is based on the following assumptions:

- Newton's law of viscosity must be obeyed by lubricant.
- The lubricant incompressible.
- The inertia forces of oil film are negligible.
- Lubricant should be of constant viscosity.

- The effect of curvature of the film with respect to film thickness is neglected.

Assumption is already made that the film is so thin that the pressure is constant across the film thickness.

- Lubricant should be constantly supplied.
- There is no slip at the boundary.
- The shaft and bearing are rigid.

An infinitesimally small element having dimensions dx , dy and dz is considered in the analysis.

u and v are the velocities in x and y direction. τ_x is the shear stress along the x direction while p is the fluid film pressure.

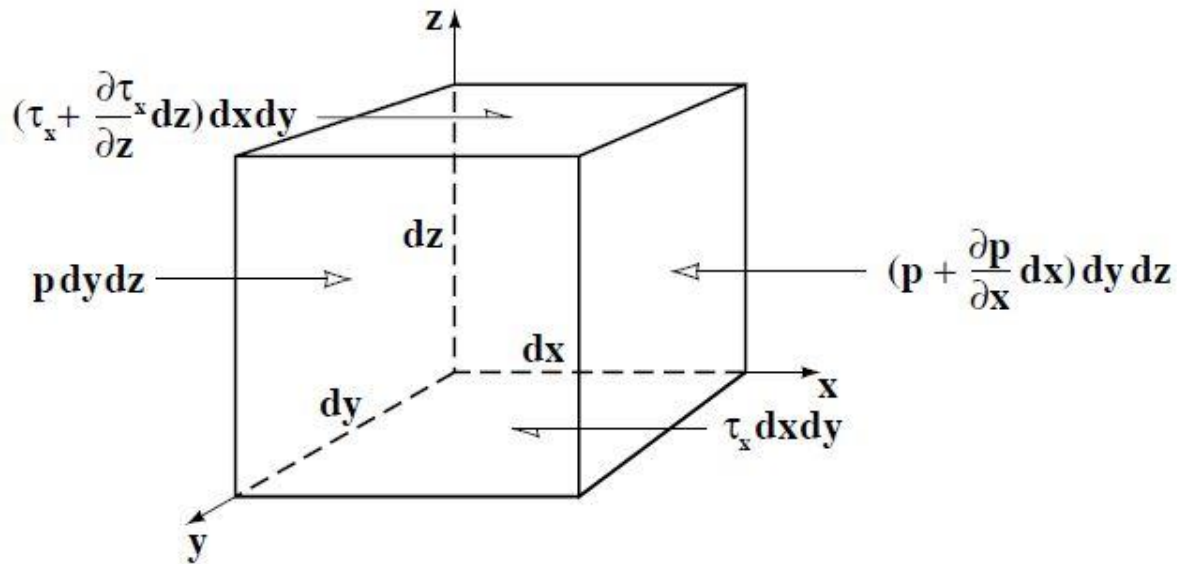


Fig 3.1 Infinitesimal element in Equilibrium (www.potto.org)

On balancing the forces in the x -direction, we get

$$p dy dz - \left(p + \frac{\partial p}{\partial x} dx \right) dy dz - \tau_x dx dz + \left(\tau_x + \frac{\partial \tau_x}{\partial y} dy \right) dx dz = 0 \quad (3.1)$$

$$\Rightarrow \frac{\partial p}{\partial x} = \frac{\partial \tau_x}{\partial z} \quad (\because \text{Assuming } dx, dy, dz \neq 0) \quad (3.2)$$

From Newton's law of viscosity we know that

$$\tau_x = \mu \frac{du}{dz} \quad (3.3)$$

Substituting (3.3) into (3.2) we get

$$\frac{\partial P}{\partial x} = \mu \frac{\partial^2 u}{\partial z^2} \quad (3.4)$$

Similarly balancing the forces in the y- direction, final result will be

$$\frac{\partial P}{\partial y} = \mu \frac{\partial^2 v}{\partial z^2} \quad (3.5)$$

And in z- direction (assumption)

$$\frac{\partial p}{\partial z} = 0 \quad (3.6)$$

As mention in the above assumption that, there is no slip or velocity at the boundary of the wedge, the boundary conditions are:

$$u = U_2 \quad \text{at} \quad z = 0$$

$$u = U_1 \quad \text{at} \quad z = h$$

Integrating twice the equation (3.4) and (3.5) we will find he value of velocities “u and v” as

$$u = \left(\frac{z^2 - zh}{2\mu} \right) \frac{\partial p}{\partial x} + (U_1 - U_2) \frac{z}{h} + U_2 \quad (3.7)$$

$$v = \left(\frac{z^2 - zh}{2\mu} \right) \frac{\partial p}{\partial y} + (V_1 - V_2) \frac{z}{h} + V_2 \quad (3.8)$$

Taking the equation of continuity

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho u) + \frac{\partial}{\partial y}(\rho v) + \frac{\partial}{\partial z}(\rho w) = 0 \quad (3.9)$$

Integrating the above equation from 0 to h

$$\int_0^h \frac{\partial \rho}{\partial t} dz + \int_0^h \frac{\partial}{\partial x}(\rho u) dz + \int_0^h \frac{\partial}{\partial y}(\rho v) dz + \int_0^h \frac{\partial}{\partial z}(\rho w) dz = 0 \quad (3.10)$$

Applying Leibnitz integration rule i.e.

$$\int_{\alpha}^{\beta} \frac{\partial u}{\partial x} dz = \frac{\partial}{\partial x} \int_{\alpha}^{\beta} u dz - u(\beta, x) \frac{d\beta}{dx} + u(\alpha, x) \frac{d\alpha}{dx}$$

Putting the value of 'u' and 'v' from equation (3.7) and (3.8) and applying at equation (3.10) a general expression for Reynolds' equation can be produced.

$$\frac{\partial}{\partial x} \left(\frac{\rho h^3}{12\mu} \frac{\partial p}{\partial x} \right) + \frac{\partial}{\partial y} \left(\frac{\rho h^3}{12\mu} \frac{\partial p}{\partial y} \right) = \frac{1}{2} \frac{\partial}{\partial x} [\rho (U_1 - U_2) h] - \rho U_2 \frac{\partial h}{\partial x} + \rho (w_2 - w_1) + h \frac{\partial \rho}{\partial t} \quad (3.11)$$

To simplify the equation further suitable assumptions are made such as neglecting the stretch, and density wedge terms. Now, Reynolds' equation becomes

$$\frac{\partial}{\partial x} \left(\frac{\rho h^3}{12\mu} \frac{\partial p}{\partial x} \right) + \frac{\partial}{\partial y} \left(\frac{\rho h^3}{12\mu} \frac{\partial p}{\partial y} \right) = \frac{1}{2} U \frac{\partial h}{\partial x} + \frac{\partial h}{\partial t} \quad (3.12)$$

Where $U = U_2 - U_1$

In polar form

$$\frac{1}{R} \frac{\partial}{\partial \theta} \left(\frac{h^3}{12\mu} \frac{1}{R} \frac{\partial P}{\partial \theta} \right) + \frac{\partial}{\partial Z} \left(\frac{h^3}{12\mu} \frac{\partial P}{\partial Z} \right) = \frac{1}{2} \frac{U}{R} \frac{\partial h}{\partial \theta} + \frac{\partial h}{\partial t} \quad (3.13)$$

The above Reynolds equation is valid for laminar fluid flow and ideal no-slip boundary condition. For finding out the pressure distribution and easy calculation we need to change the equation (3.13) in non-dimensional form. The non-dimensional form of the Reynolds' equation is given by

$$\frac{\partial}{\partial \theta} \left(\bar{h}^3 \frac{\partial \bar{P}}{\partial \theta} \right) + \left(\frac{D}{L} \right)^2 \bar{h}^3 \frac{\partial^2 \bar{P}}{\partial Z^2} = \frac{\partial \bar{h}}{\partial \theta} \quad (3.14)$$

Differentiating the above equation

$$\bar{h}^3 \frac{\partial^2 \bar{P}}{\partial \theta^2} + \frac{\partial \bar{P}}{\partial \theta} \frac{\partial \bar{h}^3}{\partial \theta} + \left(\frac{D}{L} \right)^2 \bar{h}^3 \frac{\partial^2 \bar{P}}{\partial Z^2} = \frac{\partial \bar{h}}{\partial \theta} \quad (3.15)$$

We know that

$$\bar{h} = 1 + \varepsilon \cos \theta \quad (\text{For normal journal bearing})$$

$$\bar{h} = 1 + \varepsilon \cos \theta + \bar{h}_b \quad (\text{For spherical textured journal bearing})$$

Putting the above value in equation (3.15), we get

$$\begin{aligned} \bar{h}^3 \frac{\partial^2 \bar{P}}{\partial \theta^2} - 3\varepsilon \sin \theta \frac{\partial \bar{P}}{\partial \theta} + \left(\frac{D}{L} \right)^2 \bar{h}^3 \frac{\partial^2 \bar{P}}{\partial Z^2} &= -\varepsilon \sin \theta \\ \Rightarrow \frac{\partial^2 \bar{P}}{\partial \theta^2} - \frac{3}{\bar{h}} \varepsilon \sin \theta \frac{\partial \bar{P}}{\partial \theta} + \left(\frac{D}{L} \right)^2 \bar{h}^3 \frac{\partial^2 \bar{P}}{\partial Z^2} &= -\frac{\varepsilon}{\bar{h}^3} \sin \theta \end{aligned} \quad (3.16)$$

Using central FDM

$$\begin{aligned}
\frac{\partial \bar{p}}{\partial \theta} &= \frac{\bar{p}_{i+1,j} - \bar{p}_{i-1,j}}{2\Delta\theta} \\
\frac{\partial^2 \bar{p}}{\partial \theta^2} &= \frac{\bar{p}_{i+1,j} + \bar{p}_{i-1,j} - 2\bar{p}_{i,j}}{\Delta\theta^2} \\
\frac{\partial \bar{p}}{\partial \bar{Z}} &= \frac{\bar{p}_{i,j+1} - \bar{p}_{i,j-1}}{2\Delta\bar{Z}} \\
\frac{\partial^2 \bar{p}}{\partial \bar{Z}^2} &= \frac{\bar{p}_{i,j+1} + \bar{p}_{i,j-1} - 2\bar{p}_{i,j}}{\Delta\bar{Z}^2}
\end{aligned}$$

Putting the above value in equation (3.16), equation becomes

$$\frac{\bar{p}_{i+1,j} + \bar{p}_{i-1,j} - 2\bar{p}_{i,j}}{\Delta\theta^2} - \frac{3}{1 + \varepsilon \cos \theta} \varepsilon \sin \theta \frac{\bar{p}_{i+1,j} - \bar{p}_{i-1,j}}{2\Delta\theta} + \left(\frac{D}{L}\right)^2 \frac{\bar{p}_{i,j+1} + \bar{p}_{i,j-1} - 2\bar{p}_{i,j}}{\Delta\bar{Z}^2} = -\frac{\varepsilon}{\bar{h}^3} \sin \theta$$

(3.17)

By solving equation (3.17) we can find out the value of $\bar{P}_{i,j}$ is

$$\bar{P}_{i,j} = \frac{A + B - C + D}{E}$$

Where

$$\begin{aligned}
A &= \bar{p}_{i+1,j} + \bar{p}_{i-1,j} \\
B &= \left(\frac{D}{L}\right)^2 (\bar{p}_{i,j+1} + \bar{p}_{i,j-1} - 2\bar{p}_{i,j}) \frac{\Delta\theta^2}{\Delta\bar{Z}^2} \\
C &= \frac{3}{2} \frac{\varepsilon \sin \theta}{1 + \varepsilon \cos \theta} (\bar{p}_{i+1,j} - \bar{p}_{i-1,j}) \Delta\theta \\
D &= \frac{\Delta\theta^2}{\bar{h}^3} \varepsilon \sin \theta \\
E &= 2 \left[1 + \left(\frac{D}{L}\right)^2 \frac{\Delta\theta^2}{\Delta\bar{Z}^2} \right]
\end{aligned}$$

Frictional shear force (polar form):-

$$F = \iint \tau R d\theta dz$$

$$\text{Using } \bar{\tau} = \frac{\bar{\tau} c}{\mu u}$$

Non-dimensional form is

$$= R L \int_0^L \int_0^{2\pi} \frac{u\mu}{c} \bar{\tau} d\theta \bar{dz} \quad (\text{Putting the value of } \bar{\tau})$$

$$= R L \frac{u\mu}{c} \int_0^L \int_0^{2\pi} \bar{\tau} d\theta \bar{dz}$$

$$\bar{F} = \frac{F}{R L \frac{u\mu}{c}} = \int_0^L \int_0^{2\pi} \bar{\tau} d\theta \bar{dz}$$

Load carrying capacity:-

$$W_r = - \int_0^L \int_0^{2\pi} P \cos \theta R d\theta dz$$

$$= - \frac{6u\mu R^2 L}{c^2} \int_0^L \int_0^{2\pi} \bar{P} \cos \theta d\theta \bar{dz}$$

$$W_t = \int_0^L \int_0^{2\pi} P \sin \theta R d\theta dz$$

$$= \frac{6u\mu R^2 L}{c^2} \int_0^L \int_0^{2\pi} \bar{P} \sin \theta d\theta \bar{dz}$$

$$\text{Net load } W = \sqrt{W_r^2 + W^2}$$

$$= \frac{6\mu\mu R^2 L}{c^2} \int_0^L \int_0^{2\pi} \bar{P} d\theta d\bar{z}$$

$$\bar{W} = \frac{W}{\frac{6\mu\mu R^2 L}{c^2}} = \int_0^L \int_0^{2\pi} \bar{P} d\theta d\bar{z}$$

Friction parameter :-

The force known as friction can be defined as resistance suffered by the body in relative motion by moving against each other. Coefficient of friction is defined as the ratio between the frictional force and normal load and is generally denoted by μ or η . Coefficient of friction is the characteristics of tribological system but it is not a material characteristics. The term coefficient of friction is usually used in science and engineering

We can find **friction parameter** by :-

$$\mu = \bar{F} / \bar{W}$$

Where

μ = coefficient of friction

F = frictional force

W = load carrying capacity of journal

Flow rate (Q):-

The high heat intensity developed inside the bearings of modern machines, particularly the crank pin bearing part of engine definitely requires the correct control of heat flows. Generally the heat

removal is directly proportional to the amount of lubricant flowing inside it, lubricant flow rate is one of major factors which controls the temperature inside the bearing and consequently its reliability and life.

We can find flow rate by -:

$$\bar{q} = - \int_0^L \int_0^{2\pi} \left[(\bar{h}^3) * \left(\frac{\partial \bar{p}}{\partial \bar{z}} \right) (\partial \theta) \right] / 4$$

Damping and stiffness coefficient-:

In modern era the need of increasing speed yet reliable operation of rotating machinery is increasing day by day and to achieve this we have to predict the dynamic response and stability of a rotor bearing system in an accurate way. Generally rotating machinery is supported by a one or more journal which plays a vital role in system, since it permits the relative motion between stationary and moving parts.

There are generally two types of bearings -:

1-: Fluid-film bearing

2-: Rolling-element bearing

Most of the machine of our era uses fluid-film bearing; it has large number of applications which varies from small electric motors to automobile and aircraft piston engines to large steam engines for electric power generation. In fluid-film bearing fluid-film is a thin film that separates the bearing to keep in contact with the journal, it acts like a spring and provides damping. Stiffness and damping properties of the fluid-film bearing alter the critical speed and out of balance response of rotor.

Non dimensional damping and stiffness coefficient can be calculated by:

$$\bar{k}_{\theta\theta} = \int_0^L \int_0^{2\pi} \bar{p}_x (\cos \theta) d\theta d\bar{z} \quad (3.18)$$

$$\bar{k}_{\theta z} = \int_0^L \int_0^{2\pi} \bar{p}_y (\cos \theta) d\theta d\bar{z} \quad (3.19)$$

$$\bar{k}_{z\theta} = \int_0^L \int_0^{2\pi} \bar{p}_x (\sin \theta) d\theta d\bar{z} \quad (3.20)$$

$$\bar{k}_{zz} = \int_0^L \int_0^{2\pi} \bar{p}_y (\sin \theta) d\theta d\bar{z} \quad (3.21)$$

$$\bar{b}_{\theta\theta} = \int_0^L \int_0^{2\pi} \bar{p}_x^* (\cos \theta) d\theta d\bar{z} \quad (3.22)$$

$$\bar{b}_{\theta z} = \int_0^L \int_0^{2\pi} \bar{p}_y^* (\cos \theta) d\theta d\bar{z} \quad (3.23)$$

$$\bar{b}_{z\theta} = \int_0^L \int_0^{2\pi} \bar{p}_x^* (\sin \theta) d\theta d\bar{z} \quad (3.24)$$

$$\bar{b}_{zz} = \int_0^L \int_0^{2\pi} \bar{p}_y^* (\sin \theta) d\theta d\bar{z} \quad (3.25)$$

Where

$\bar{k}_{\theta\theta}$, $\bar{k}_{z\theta}$, $\bar{k}_{\theta z}$, \bar{k}_{zz} are the non-dimensional stiffness in different directions and $\bar{b}_{\theta\theta}$, $\bar{b}_{z\theta}$, $\bar{b}_{\theta z}$, \bar{b}_{zz} are the non-dimensional damping coefficients in different directions which are calculated by infinitely perturbation method [50].

The governing equation for the pressure distribution in hydro dynamic journal bearing film is:-

$$\frac{1}{R^2} \frac{\partial}{\partial \theta} \left(\frac{h^3}{12\eta} \frac{\partial p}{\partial \theta} \right) + \frac{\partial}{\partial z} \left(\frac{h^3}{12\eta} \frac{\partial p}{\partial z} \right) = \frac{1}{2} \omega \frac{\partial h}{\partial \theta} + \frac{\partial h}{\partial t} \quad (3.26)$$

Now by changing the governing pressure distribution equation in to non-dimensional form we get:-

$$\frac{\partial}{\partial \theta} \left(\bar{h}^3 \frac{\partial \bar{p}}{\partial \theta} \right) + \frac{R^2}{L^2} \bar{h}^3 \frac{\partial^2 \bar{p}}{\partial \bar{z}^2} = \frac{\partial \bar{h}}{\partial \theta} \quad (3.27)$$

Under dynamic conditions the journal center motion is described by the amplitudes Δx and Δy , measured from the static equilibrium position, such that the film thickness becomes:-

$$h = h_0 + \Delta x \cos \theta + \Delta y \sin \theta \quad (3.28)$$

$$p = p_0 + p_x \Delta x + p_y \Delta y + p_x^\bullet \Delta x^\bullet + p_y^\bullet \Delta y^\bullet \quad (3.29)$$

By substituting equation 3.29 and 3.28 in equation 3.26 we get:-

$$\left\{ \frac{1}{R^2} \frac{\partial}{\partial \theta} \left(\frac{h_0^3}{12\eta} \frac{\partial p}{\partial \theta} \right) + \frac{\partial}{\partial z} \left(\frac{h_0^3}{12\eta} \frac{\partial p}{\partial z} \right) \right\} \begin{Bmatrix} p_0 \\ p_x \\ p_y \\ p_x^\bullet \\ p_y^\bullet \end{Bmatrix} = \begin{Bmatrix} \frac{1}{2} \omega \frac{\partial h_0}{\partial \theta} \\ -\frac{1}{2} \omega [\sin \theta + 3 \frac{\cos \theta}{h_0} \frac{\partial h}{\partial \theta}] - \frac{h_0^3}{4\eta} \frac{1}{R^2} \frac{\partial p_0}{\partial \theta} \frac{\partial}{\partial \theta} \left[\frac{\cos \theta}{h_0} \right] \\ \frac{1}{2} \omega [\cos \theta - 3 \frac{\sin \theta}{h_0} \frac{\partial h}{\partial \theta}] - \frac{h_0^3}{4\eta} \frac{1}{R^2} \frac{\partial p_0}{\partial \theta} \frac{\partial}{\partial \theta} \left[\frac{\sin \theta}{h_0} \right] \\ \cos \theta \\ \sin \theta \end{Bmatrix} \quad (3.30)$$

Using above equations we will calculate p_0

$$\frac{1}{R^2} \frac{\partial}{\partial \theta} \left(\frac{h_0^3}{12\eta} \frac{\partial p_0}{\partial \theta} \right) + \frac{\partial}{\partial z} \left(\frac{h_0^3}{12\eta} \frac{\partial p_0}{\partial z} \right) = \frac{1}{2} \omega \frac{\partial h_0}{\partial \theta} \quad (3.31)$$

First of all we will non dimensionlized the above equation so that all the values which we are calculating can be used universally for all the bearings.

By multiplying equation 3.31 by $12\eta R^2$ we get-

$$\frac{\partial}{\partial \theta} \left(h_0^3 \frac{\partial p_0}{\partial \theta} \right) + R^2 h_0^3 \frac{\partial^2 p_0}{\partial z^2} = \frac{1}{2} \omega \frac{\partial h_0}{\partial \theta} * 12\eta R^2 \quad (3.32)$$

Making the above equation dimensionless by using the following,

$$\bar{h} = \frac{h}{c}, \quad \bar{p}_0 = \frac{p_0 c^2}{6\eta \omega R^2}, \quad \bar{z} = \frac{z}{L}$$

$$\frac{\partial}{\partial \theta} \left(\bar{h}_0^3 c^3 \frac{\partial \bar{p}_0}{\partial \theta} * \frac{6\eta \omega R^2}{c^2} \right) + \left(\frac{R^2}{L^2} \right) \bar{h}_0^3 c^3 \frac{\partial^2 \bar{p}_0}{\partial \bar{z}^2} \frac{6\eta \omega R^2}{c^2} = \frac{1}{2} \omega \frac{\partial \bar{h}_0}{\partial \theta} 12\eta R^2 c \quad (3.33)$$

Eliminate $6\eta \omega R^2 c$ from whole equation

$$\frac{\partial}{\partial \theta} \left(\bar{h}_0^3 \frac{\partial \bar{p}_0}{\partial \theta} \right) + \left(\frac{R^2}{L^2} \right) \bar{h}_0^3 \frac{\partial^2 \bar{p}_0}{\partial \bar{z}^2} = \frac{\partial \bar{h}_0}{\partial \theta} \quad (3.34)$$

$$\left(\bar{h}_0^3 \frac{\partial^2 \bar{p}_0}{\partial \theta^2} + \frac{\partial \bar{p}_0}{\partial \theta} \right) 3\bar{h}_0^2 \frac{\partial \bar{h}_0}{\partial \theta} + \left(\frac{R^2}{L^2} \right) \bar{h}_0^3 \frac{\partial^2 \bar{p}_0}{\partial \bar{z}^2} = \frac{\partial \bar{h}_0}{\partial \theta} \quad (3.35)$$

The above equation 3.35 is the non dimensionlized form of equation 3.31.

Now on dividing whole equation 3.35 by \bar{h}_0^3 we get:-

$$\frac{\partial^2 \bar{p}_0}{\partial \theta^2} + \left(\frac{\partial \bar{p}_0}{\partial \theta} \right) \frac{3}{\bar{h}_0} \left(\frac{\partial \bar{h}_0}{\partial \theta} \right) + \left(\frac{R^2}{L^2} \right) \frac{\partial^2 \bar{p}_0}{\partial \bar{z}^2} = \frac{1}{\bar{h}_0^3} \frac{\partial \bar{h}_0}{\partial \theta} \quad (3.36)$$

On applying finite difference method in equation 3.36 we get pressure in the form of:-

$$\bar{p}_0 = \frac{A+B+C+D}{E}$$

$$\text{Where } A = \bar{P}_{o_{i+1,j}} + \bar{P}_{o_{i-1,j}}$$

$$B = \frac{R^2}{L^2} [\bar{P}_{o_{i,j+1}} + \bar{P}_{o_{i,j-1}}] \left(\frac{\Delta\theta^2}{\Delta\bar{z}^2} \right)$$

$$C = -\frac{3}{2h_0} \varepsilon \sin \theta \Delta\theta [\bar{P}_{o_{i+1,j}} - \bar{P}_{o_{i-1,j}}]$$

$$D = \varepsilon \sin \theta \frac{\Delta\theta^2}{\bar{h}_0^3}$$

$$E = 2 \left[1 + \left(\frac{R^2}{L^2} \right) \frac{\Delta\theta^2}{\Delta\bar{z}^2} \right]$$

Similarly we can calculate \bar{p}_x , \bar{p}_y , \bar{p}_x^\bullet , \bar{p}_y^\bullet and by putting these all in equation 3.18 to 3.25 we will calculate stiffness and damping coefficients.

$$\bar{p}_x = \frac{A_1+B_1+C_1+D_1}{E_1}$$

$$A_1 = \bar{P}_{x_{i+1,j}} + \bar{P}_{x_{i-1,j}}$$

$$B_1 = \frac{R^2}{L^2} [\bar{P}_{x_{i,j+1}} + \bar{P}_{x_{i,j-1}}] \left(\frac{\Delta\theta^2}{\Delta\bar{z}^2} \right)$$

$$C_1 = -\frac{3}{2h_0} \varepsilon \sin \theta \Delta\theta [\bar{P}_{x_{i+1,j}} - \bar{P}_{x_{i-1,j}}]$$

$$D_1 = \frac{\Delta\theta^2}{\bar{h}_0^3} \left[\left\{ \sin \theta - \frac{3\varepsilon \cos \theta \sin \theta}{\bar{h}_0} \right\} + 3\bar{h}_0 (-\bar{h}_0 \sin \theta + \varepsilon \cos \theta \sin \theta) \left(\frac{\bar{P}_{o_{i+1,j}} - \bar{P}_{o_{i-1,j}}}{2\Delta\theta} \right) \right]$$

$$E_1 = 2 \left[1 + \left(\frac{R^2}{L^2} \right) \frac{\Delta \theta^2}{\Delta \bar{z}^2} \right]$$

$$\bar{p}_y = \frac{A_2 + B_2 + C_2 + D_2}{E_2}$$

$$A_2 = \bar{P}_{y_{i+1,j}} + \bar{P}_{y_{i-1,j}}$$

$$B_2 = \frac{R^2}{L^2} [\bar{P}_{y_{i,j+1}} + \bar{P}_{y_{i,j-1}}] \left(\frac{\Delta \theta^2}{\Delta \bar{z}^2} \right)$$

$$C_2 = -\frac{3}{2h_0} \varepsilon \sin \theta \Delta \theta [\bar{P}_{y_{i+1,j}} - \bar{P}_{y_{i-1,j}}]$$

$$D_2 = - \left(\frac{\Delta \theta^2}{\bar{h}_0^3} \left[\left\{ \cos \theta + \frac{3\varepsilon \sin^2 \theta}{\bar{h}_0} \right\} - 3\bar{h}_0 (\bar{h}_0 \cos \theta + \varepsilon \sin^2 \theta) \left(\frac{\bar{p}_{0_{i+1,j}} - \bar{p}_{0_{i-1,j}}}{2\Delta \theta} \right) \right] \right)$$

$$E_2 = 2 \left[1 + \left(\frac{R^2}{L^2} \right) \frac{\Delta \theta^2}{\Delta \bar{z}^2} \right]$$

Similarly for \bar{p}_x^\bullet

$$\bar{p}_x^\bullet = \frac{A_3 + B_3 + C_3 + D_3}{E_3}$$

$$A_3 = \bar{p}_{x_{i+1,j}}^\bullet + \bar{p}_{x_{i-1,j}}^\bullet$$

$$B_3 = \frac{R^2}{L^2} [\bar{p}_{x_{i,j+1}}^\bullet + \bar{p}_{x_{i,j-1}}^\bullet] \left(\frac{\Delta \theta^2}{\Delta \bar{z}^2} \right)$$

$$C_3 = -\frac{3}{2h_0} \varepsilon \sin \theta \Delta \theta [\bar{p}_{x_{i+1,j}}^\bullet - \bar{p}_{x_{i-1,j}}^\bullet]$$

$$D_3 = \left[-2 \left(\frac{\cos \theta}{h_0^3} \right) * \Delta \theta^2 \right]$$

$$E_3 = 2 \left[1 + \left(\frac{R^2}{L^2} \right) \frac{\Delta \theta^2}{\Delta \bar{z}^2} \right]$$

Similarly for \bar{p}_y^\bullet

$$\bar{p}_y^\bullet = \frac{A_4 + B_4 + C_4 + D_4}{E_4}$$

$$A_4 = \bar{p}_{y_{i+1,j}}^\bullet + \bar{p}_{y_{i-1,j}}^\bullet$$

$$B_4 = \frac{R^2}{L^2} [\bar{p}_{y_{i,j+1}}^\bullet + \bar{p}_{y_{i,j-1}}^\bullet] \left(\frac{\Delta \theta^2}{\Delta \bar{z}^2} \right)$$

$$C_4 = -\frac{3}{2h_0} \varepsilon \sin \theta \Delta \theta [\bar{p}_{y_{i+1,j}}^\bullet - \bar{p}_{y_{i-1,j}}^\bullet]$$

$$D_4 = \left[-2 * \left(\frac{\sin \theta}{h_0^3} \right) * \Delta \theta^2 \right]$$

$$E_4 = 2 \left[1 + \left(\frac{R^2}{L^2} \right) \frac{\Delta \theta^2}{\Delta \bar{z}^2} \right]$$

By putting values of $\bar{p}_x, \bar{p}_y, \bar{p}_x^\bullet, \bar{p}_y^\bullet$ in equations 3.18 to 3.25, we will get various values of stiffness and damping coefficient.

Flow Chart to calculate pressure profile and other parameters of spherical textured journal bearing

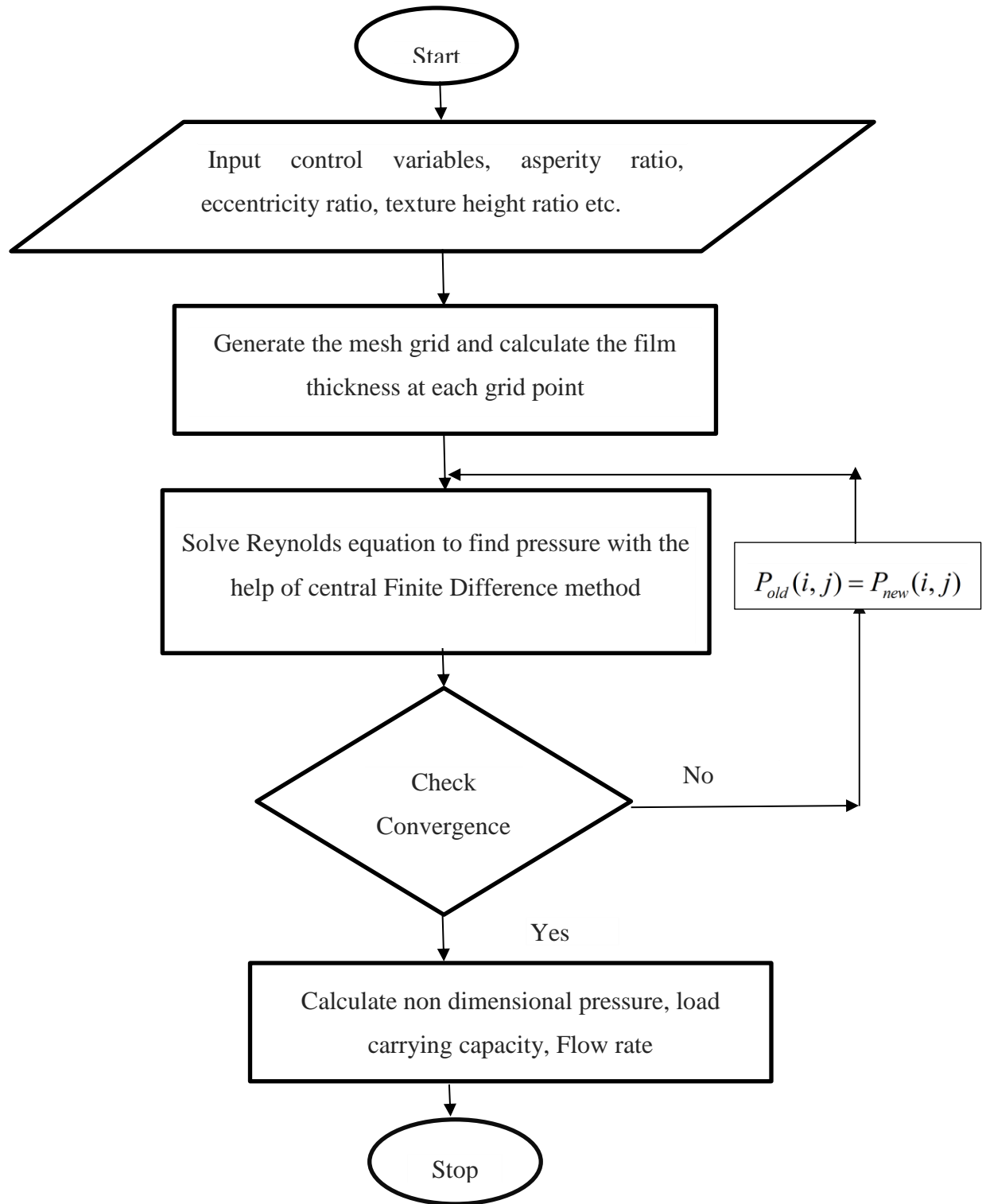


Fig. 3.2 Flow chart (static characteristics)

Flow chart to calculate stiffness and damping coefficient using infinitely perturbation method

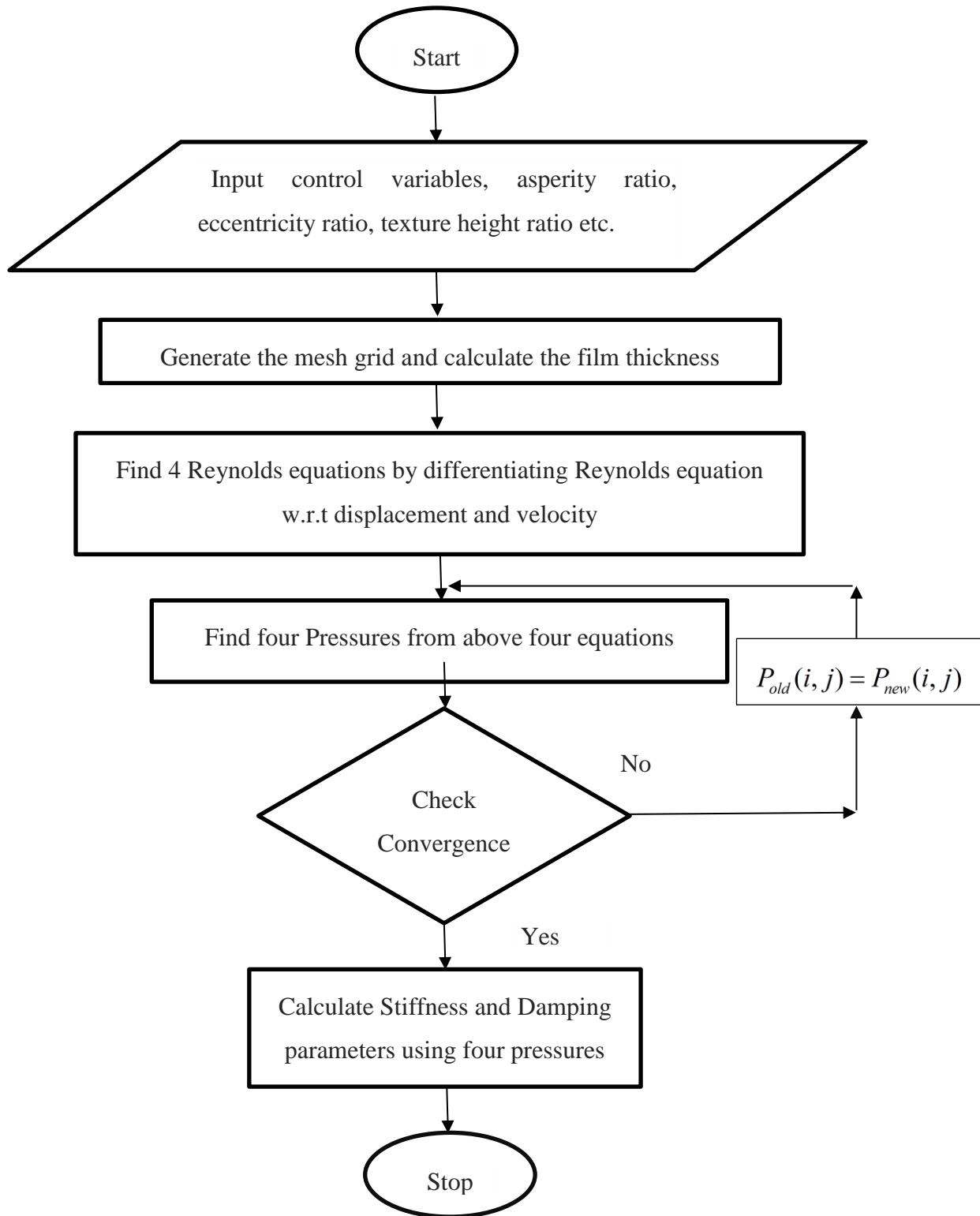


Fig.3.3 Flow chart (dynamic characteristics)

Result and Discussion

Theoretical and graphical analysis of texture is done on the hydrodynamic journal bearing to study its effect on parameters of journal bearing. In present analysis we are studying the effect of eccentricity ratio, asperity ratio and texture height on non-dimensional load, friction parameter, non-dimensional flow rate, stiffness coefficients and damping coefficients.

4.1 What is aspect ratio and texture height ratio?

For numerical analysis, the bearing surface is modeled by dividing into finite number of squares (unit cell) which are equal to number of asperities, i.e. each unit cell contain one asperity. Aspect Ratio is defined as ratio of projected area of asperities to total area of square. Similarly each asperity has some height, when the maximum height of the asperity is divided by the width of the unit cell it gives the texture height ratio. Figure 4.1 shows asperities with different aspect ratio and height ratio for better understanding of the terms.

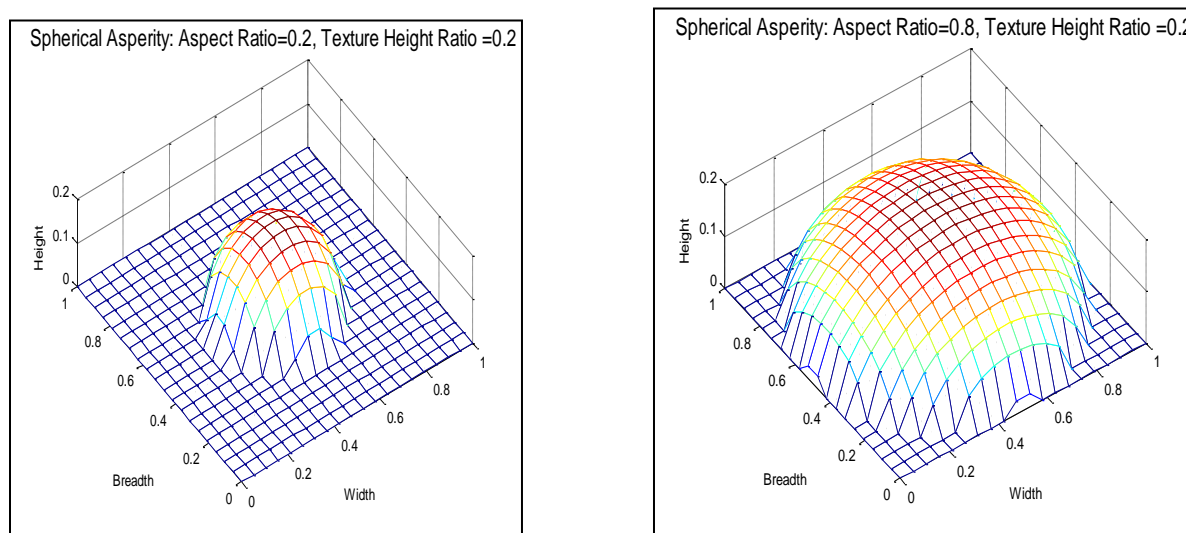


Fig.4.1 Modelling with different aspect ratio and texture height ratio

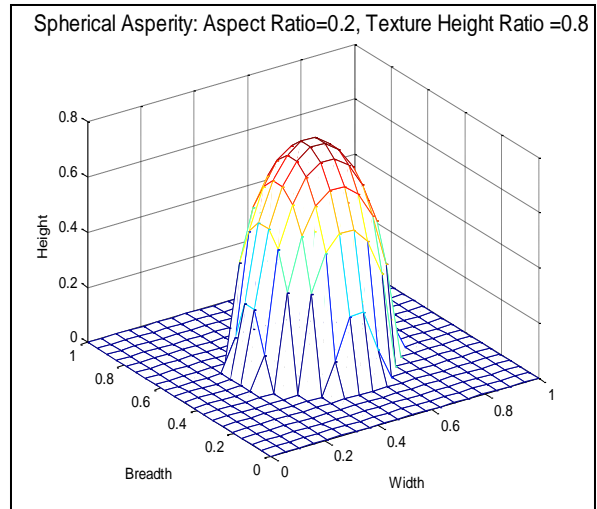
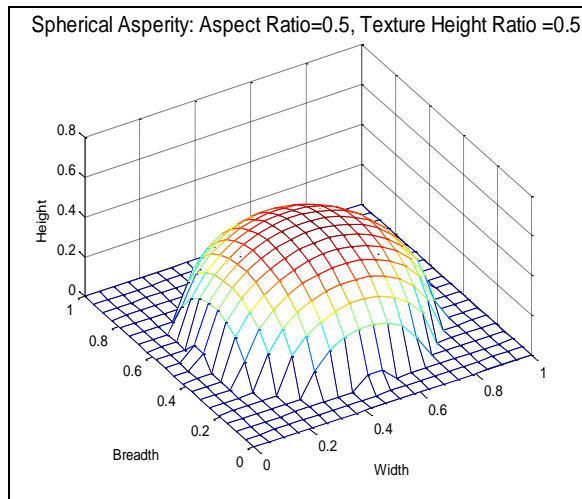


Fig.4.2 Modelling with different aspect ratio and texture height ratio

How spherical textured hydrodynamic bearing does looks alike?

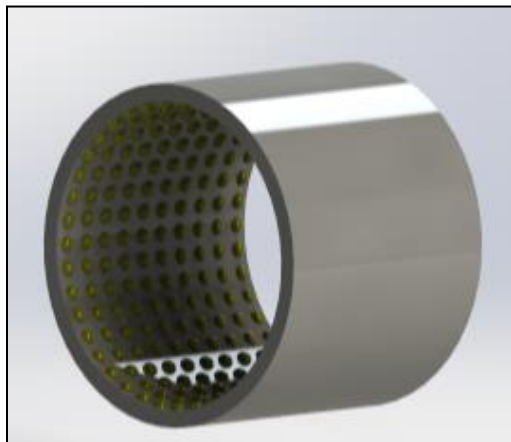


Fig.4.3 Textured bearing

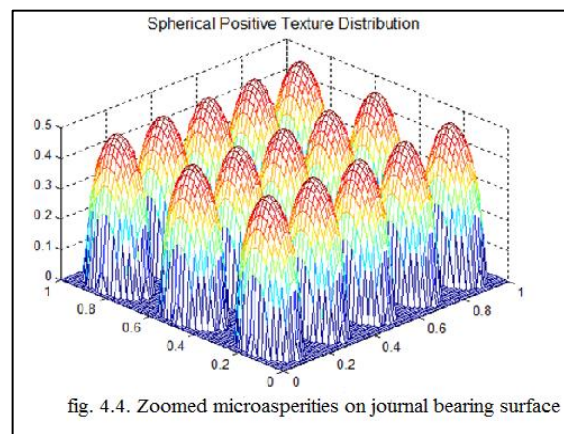


Fig 4.4 Zoomed micro asperities on journal bearing surface

Texture bearing have undulations or micro asperities or dimples or texture which are created inside the surface by bearing by some surface texturing techniques, the bearing will look like as shown in fig.4.3 . Fig.4.4 shows the zoomed view of spherical micro asperities on the surface of bearing.

4.2 Spherical texture

Spherical texture can be positive or negative as shown in fig 4.5 and fig 4.7 but we have taken positive spherical texture to analyze the effect of texture on the parameters of journal bearing.

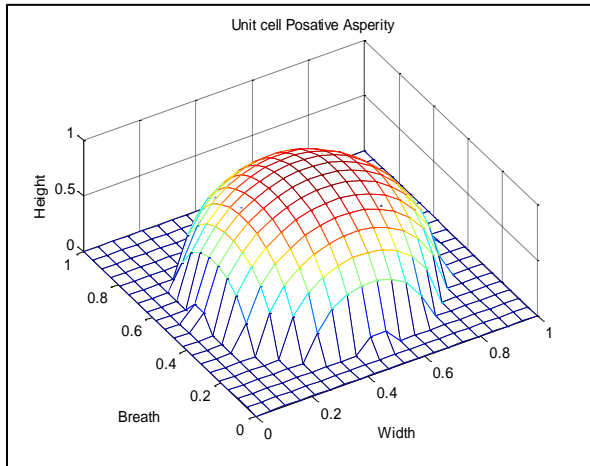


Fig 4.5 Single positive asperity

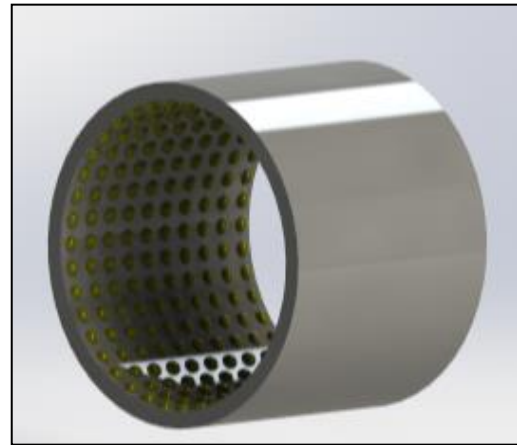


Fig 4.6 Distributed positive asperity

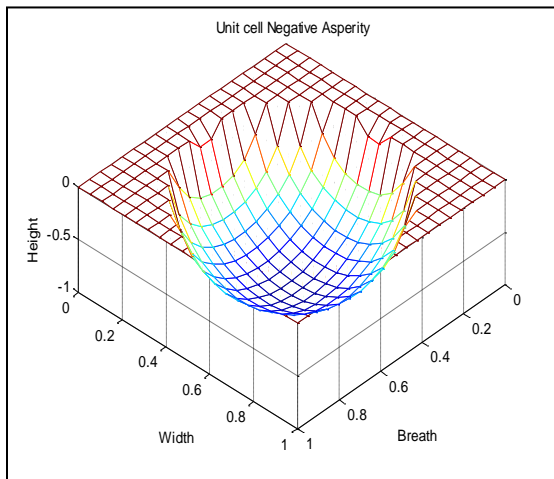


Fig 4.7 Single negative asperity

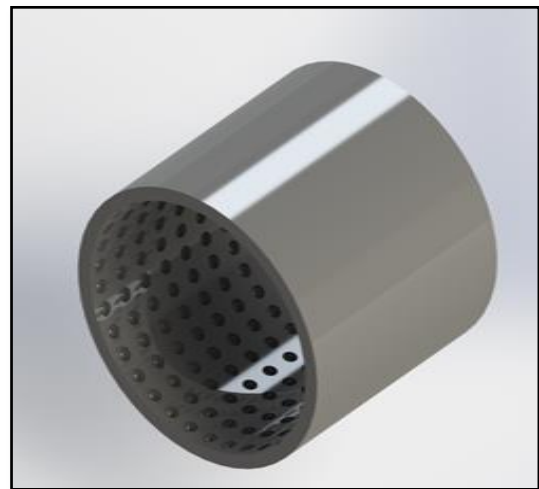


Fig 4.8 Distributed negative asperity

Above are micron scaled surface featured with spherical asperity on a surface. Model of single positive asperity and distributed positive asperity is shown in figure 4.5 and figure 4.6 whereas model of single negative asperity and distributed negative asperity is shown in figure 4.7 and figure 4.8.

Asperities can be of many types such as square, rectangular, elliptical, cylindrical, spherical etc. also they can be positive or negative but from all these types of asperities we are taking only spherical asperity and its effects on the parameters of hydrodynamic journal bearing because from all these types of asperities only spherical asperity will give optimum result.

4.3 Comparison between pressure profile of normal and negative spherical textured journal bearing on varying eccentricity ratio

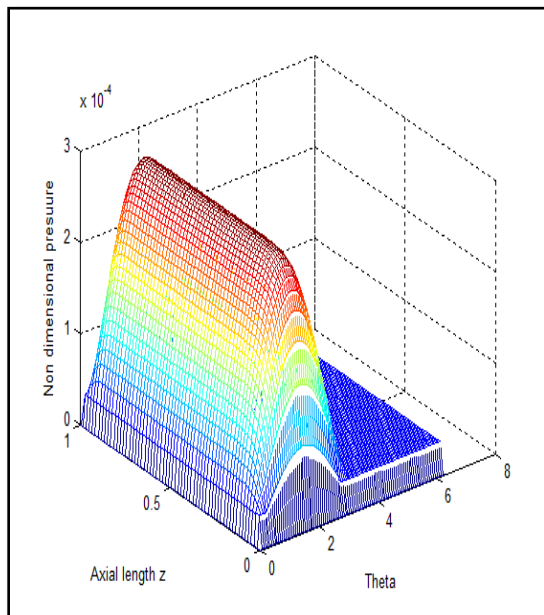


Fig. 4.9 Pressure profile of normal journal bearing

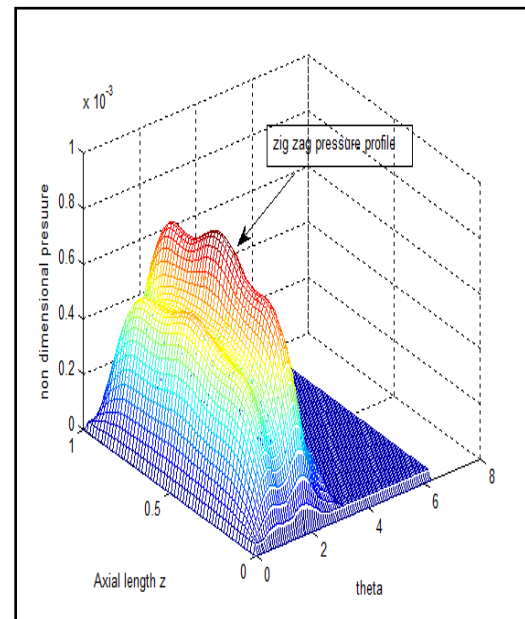


Fig. 4.10 Pressure profile of spherical textured journal bearing

Figure 4.9 and 4.10 shows the pressure profile of both normal journal bearing and spherical textured journal bearing in theta-z plane. We have taken a non-dimensional pressure and from this non dimensionlized pressure we can calculate the actual pressure by multiplying it with certain terms.

From figure 4.9 and 4.10 we see that the pressure profile for normal journal bearing is smooth as compared to the pressure profile of spherical textured journal bearing. Spherical textured hydrodynamic journal bearing has a textured bearing in which spherical shape texture or undulation or asperity or micro dimples are created on the surface of bearing due to which there will be more convergence between the shaft and bearing and due to this more convergence, more pressure will be developed inside the spherical textured hydrodynamic bearing as compared to the normal journal bearing. From figure we can see that spherical texture journal bearing has more pressure as compared to normal journal bearing and also the profile of spherical textured journal bearing is in a zig-zag manner while pressure profile of normal journal bearing is smooth as shown by arrow in figure 4.10.

4.4 Effect on non-dimensional load carrying capacity of both spherical textured and normal journal bearing on varying eccentricity ratio

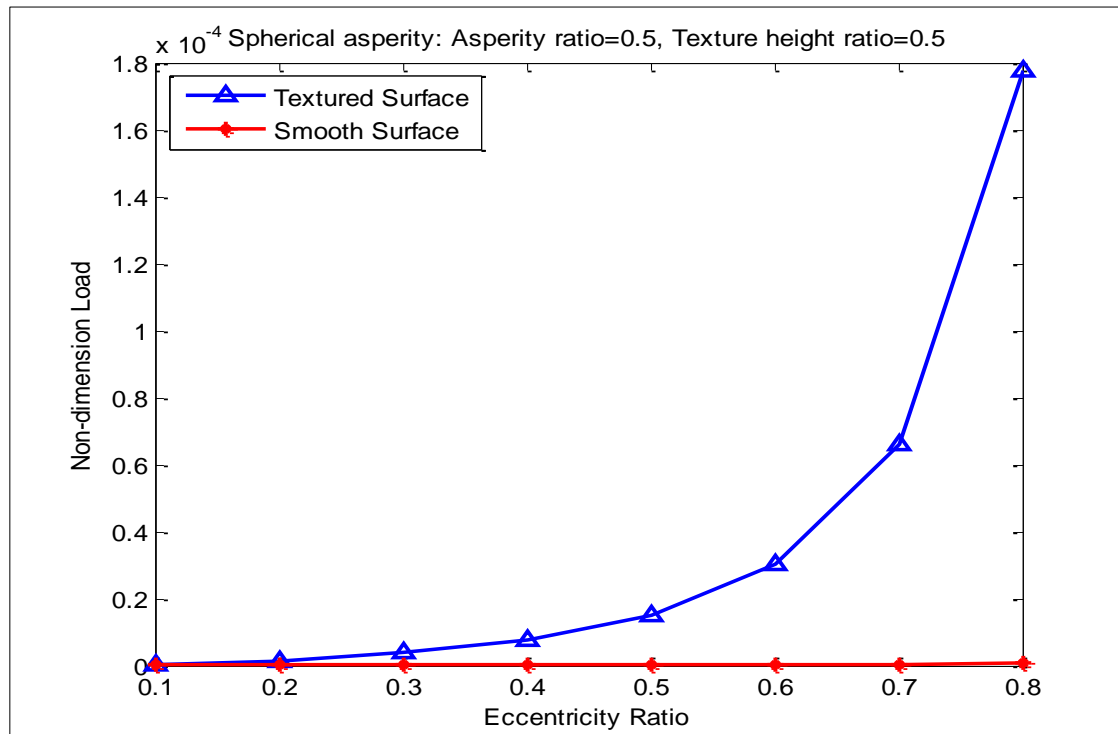


Fig 4.11 Eccentricity ratio vs Non dimensional load

The above graph is in between eccentricity ratio vs non dimensional load for negative spherical asperity. We have taken eccentricity ratio as abscissa because radial clearance is in our control, we can increase or decrease eccentricity ratio by varying radial clearance. Eccentricity ratio is the ratio of eccentricity to the radial clearance. From fig 4.11 we get that on increasing eccentricity ratio load carrying capacity of both normal and textured bearing increases, but in latter case it increases more than normal journal bearing because in textured journal bearing more pressure will be there due to the more convergence which happen due to the texture formation on the surface of bearing and due to this more pressure it can carry large amount of

load, so by varying eccentricity ratio from 0.1 to 0.8 we see that after reaching eccentricity ratio 0.6 there is an sudden increment in load carrying capacity of texture journal bearing.

4.5 Effect on non-dimensional friction parameter of both normal and spherical textured journal bearing on varying eccentricity ratio

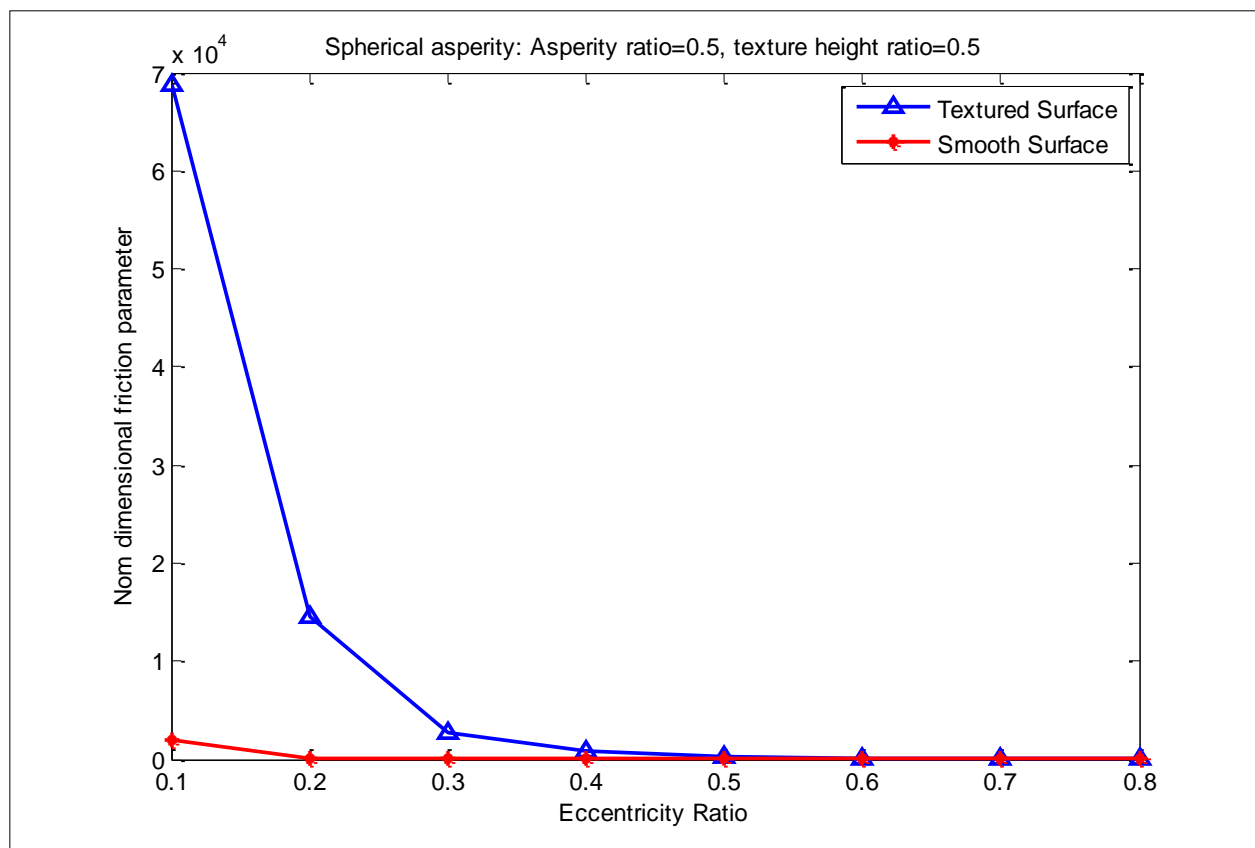


Fig. 4.12 Eccentricity ratio vs friction parameter

In industries friction is the term which is generally used now and then, friction is directly proportional to power, more the friction, more will be the wastage of power. In order to minimize the wastage of power we have to control friction. In above figure blue color is for textured journal bearing and red color is for normal journal bearing. As we can see from the

figure that during starting, non-dimensional coefficient of friction parameter is more in textured journal bearing than normal journal bearing, but as soon as we reach the eccentricity ratio i.e. 0.6 it becomes almost equal or less after that as shown in figure 4.12.

4.6 Effect of surface texture on lubricant flow rate on both spherical textured and normal journal bearing on varying eccentricity ratio

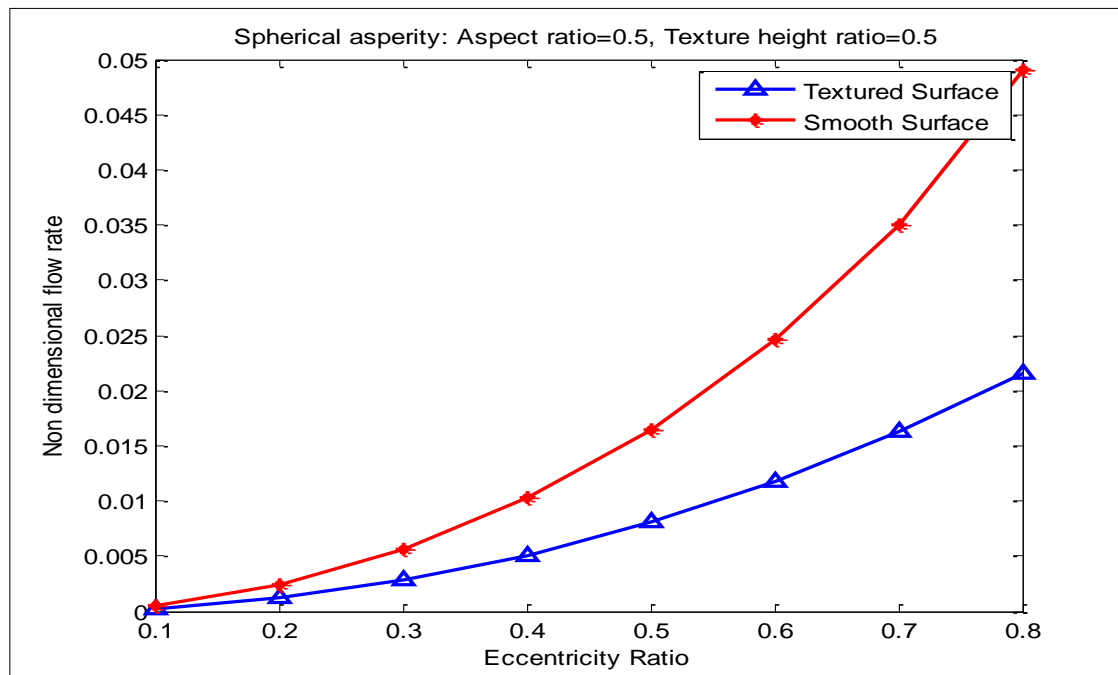


Fig 4.13 Eccentricity ratio vs Non dimensional flow rate

To control friction in journal bearing, lubrication is provided inside the journal bearing so that we can control power loss by reducing friction. Figure 4.13 shows a comparison of flow rate of lubricant inside the normal journal bearing and inside the spherically textured journal bearing. From figure we can see that lubricant flow rate in textured journal bearing is less than that of normal journal bearing, it is because, in spherically textured journal bearing, texture behaves as a

reservoir which stores lubricant and provide it during starving condition, so there is less wastage of lubricant in textured journal bearing as compared to normal bearing.

4.7 Variation of various stiffness coefficients on varying eccentricity ratio for normal journal bearing

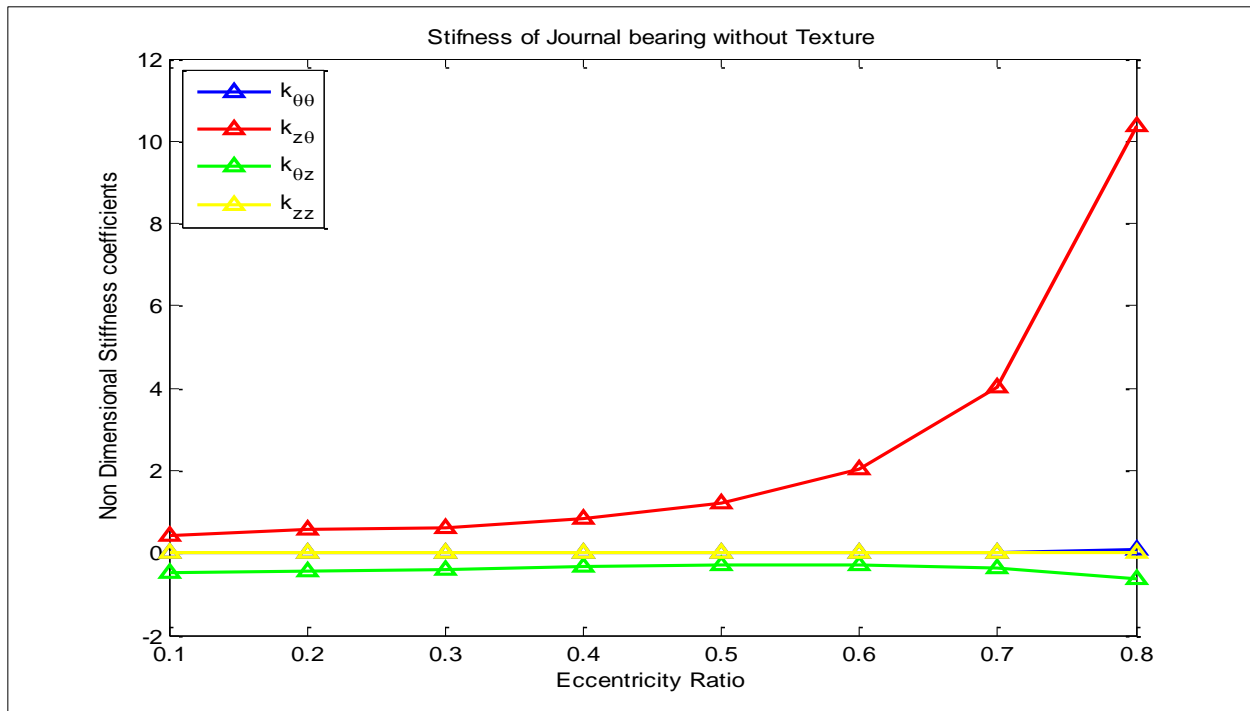


Fig 4.14 Eccentricity ratio vs Non dimensional stiffness coefficient (normal bearing)

Above figure shows the variation of various non dimensional stiffness coefficients on varying eccentricity ratio for normal journal bearing. $k_{\theta\theta}$ and k_{zz} lines on graph have constant slope, whereas the slope of $k_{z\theta}$, $k_{\theta z}$ line on graph is positive and increasing. Thus we can conclude from the above figure that on increasing eccentricity ratio stiffness coefficient $k_{z\theta}$ and $k_{\theta z}$ increases while the other two stiffness coefficients remain constant.

4.8 Variation of various stiffness coefficients on varying eccentricity ratio for negative spherical textured journal bearing

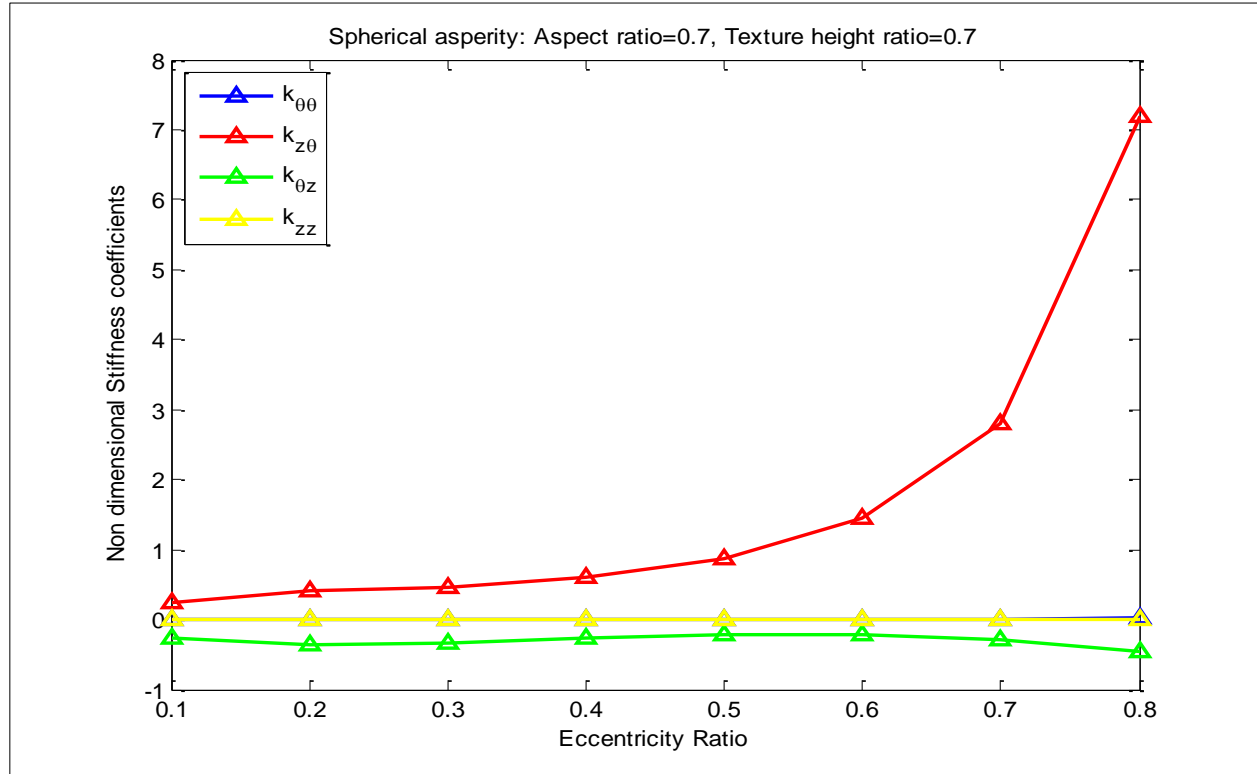


Fig 4.15 Eccentricity ratio vs Non dimensional stiffness coefficients

Fluid film stiffness coefficients are affected by the behavior of lubricant, from figure 4.14 we can see that on increasing eccentricity ratio, stiffness coefficient of the spherically textured journal bearing $k_{z\theta}$ and $k_{\theta z}$ increases whereas both the other stiffness coefficients have no effect or very small effect. On comparing figure 4.14 and 4.15 we get that stiffness coefficient of spherically textured journal bearing is slightly less than that of normal journal bearing. The role of stiffness is to provide the journal a restoring force that pushes back journal towards its steady state equilibrium condition.

4.9 Variation of various damping coefficient on varying eccentricity ratio for normal journal bearing

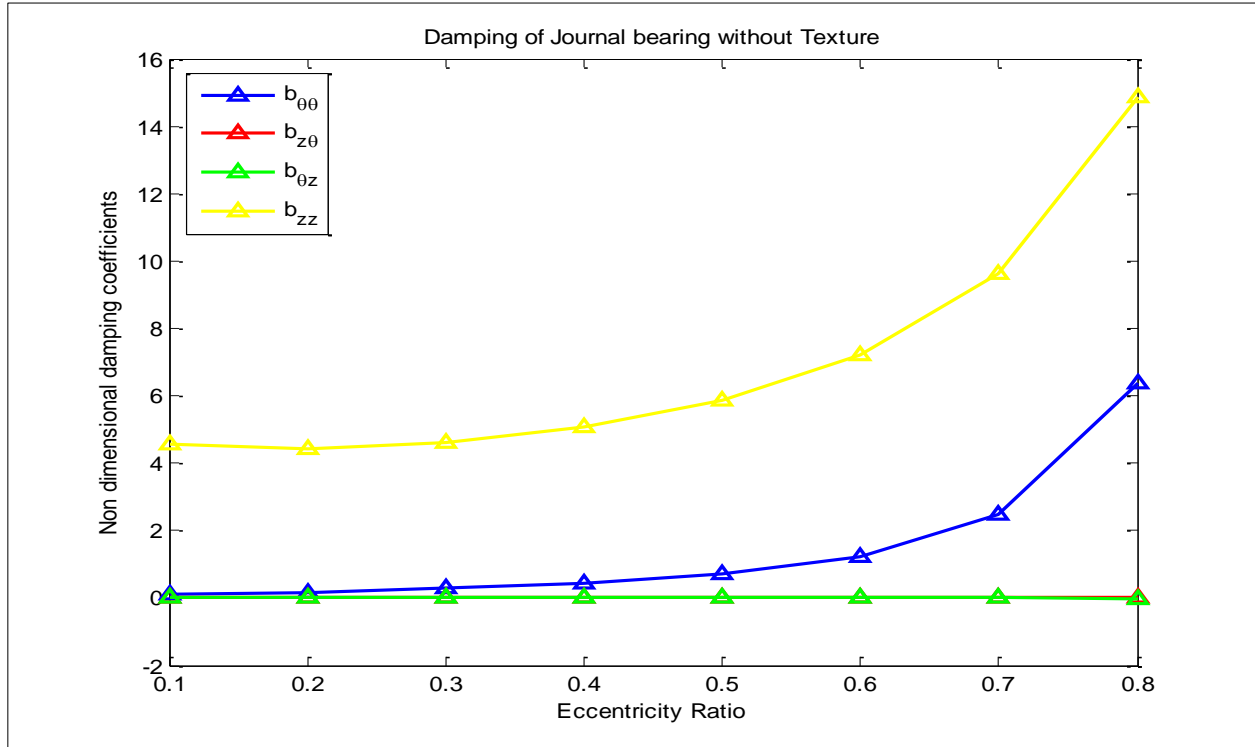


Fig 4.16 Eccentricity ratio vs Non dimensional damping coefficient (normal bearing)

Above figure shows the variation of damping coefficients of normal journal bearing on varying eccentricity ratio. From above figure we can see that damping coefficient $b_{\theta z}$ and $b_{z\theta}$ are almost equal while damping coefficient $b_{\theta\theta}$ and b_{zz} increases on increasing eccentricity ratio, later one initially has larger value than the former one as shown in figure. Slope of both $b_{\theta\theta}$ and b_{zz} are positive and increasing.

Damping is generally provided for following purpose:

- In order to pass through critical speeds.
- To suppress instabilities and sub synchronous vibration.

- To reduce noise.
- Able to withstand shock loads
- Reduced transmitted vibration.

4.10 Variation of various damping coefficient of negative spherical textured journal bearing on varying eccentricity ratio

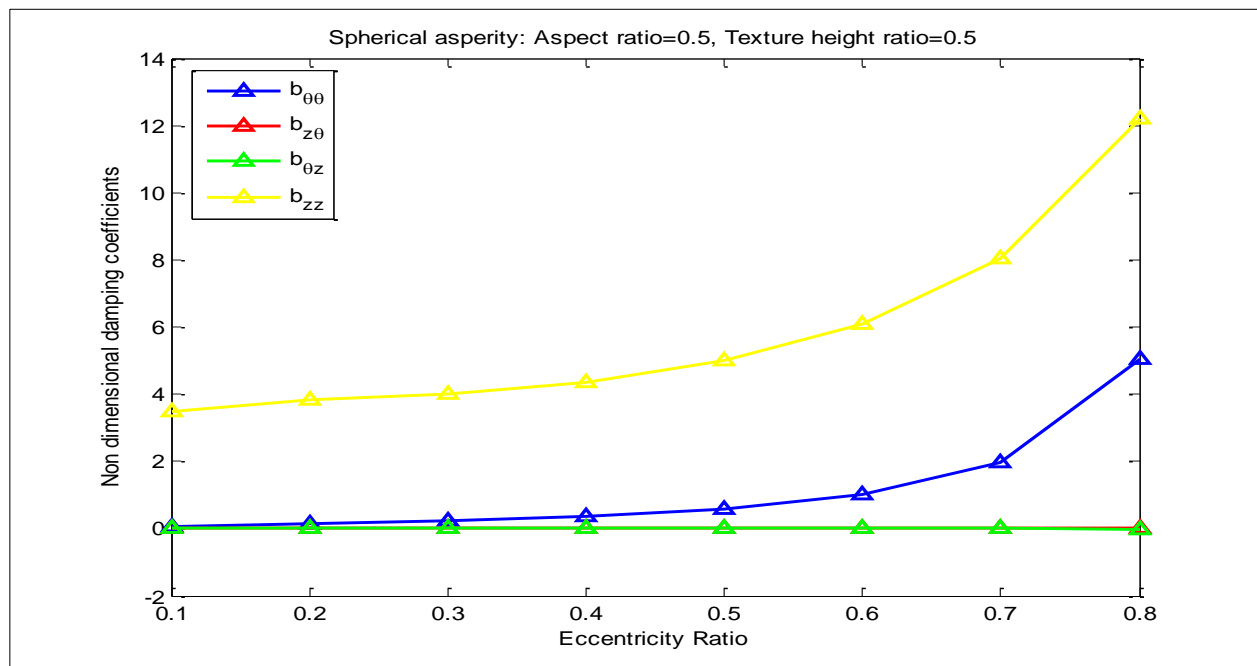


Fig. 4.17 Eccentricity ratio vs Non dimensional damping coefficients (textured bearing)

From above figure we can see that for a spherically textured hydrodynamic journal bearing, on increasing the eccentricity ratio of spherically textured journal bearing, damping coefficient $b_{z\theta}$ and $b_{\theta z}$ remains unchanged whereas damping coefficients $b_{\theta\theta}$ and b_{zz} increases with it, the later one has initially higher value than the former one. On comparing figure 4.16 and 4.17 we get that damping coefficients of textured bearing is slightly less than the normal bearing.

4.11 Variation of load for different number of textures on varying asperity ratio for spherically textured journal bearing

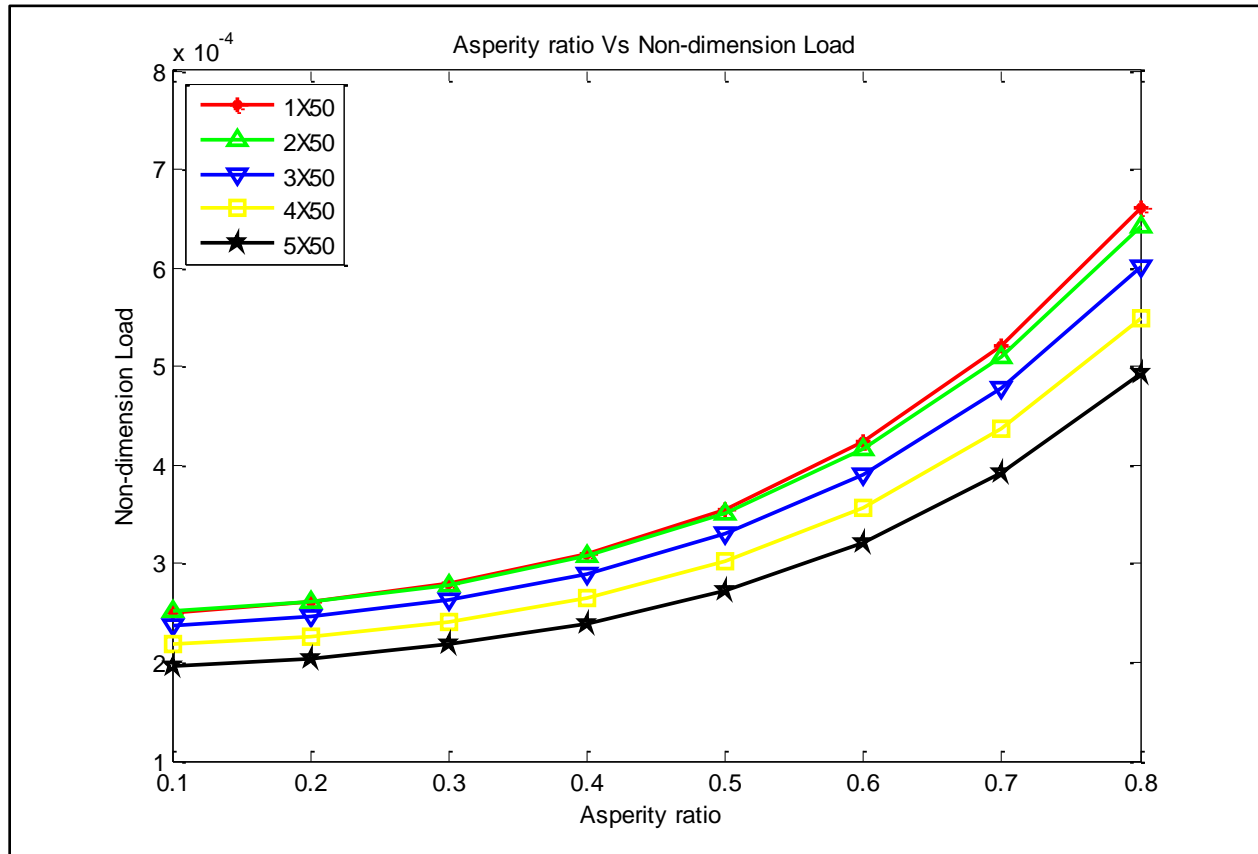


Fig 4.18 Asperity ratio vs Non dimensional load

From above figure we can see that for spherically textured hydrodynamic journal bearing, on varying asperity ratio (ratio of area of dimple to the area of cell) for different number of textures how non dimensional load is varying. Each line of different colors indicates the number of textures used to calculate non dimension load by varying asperity ratio for e.g. black line is for 5X50 number of textures, and on varying asperity ratio, non-dimension load also increase therefore the slope of black line is positive, from figure we observe that the surface having less number of textures have slightly high load carrying capacity than others because as soon as we

cross the limit of increasing the number of textures , the surface instead of becoming rough will become smooth.

4.12 Variation of coefficient of friction parameter for different number of textures on varying asperity ratio for spherically textured journal bearing

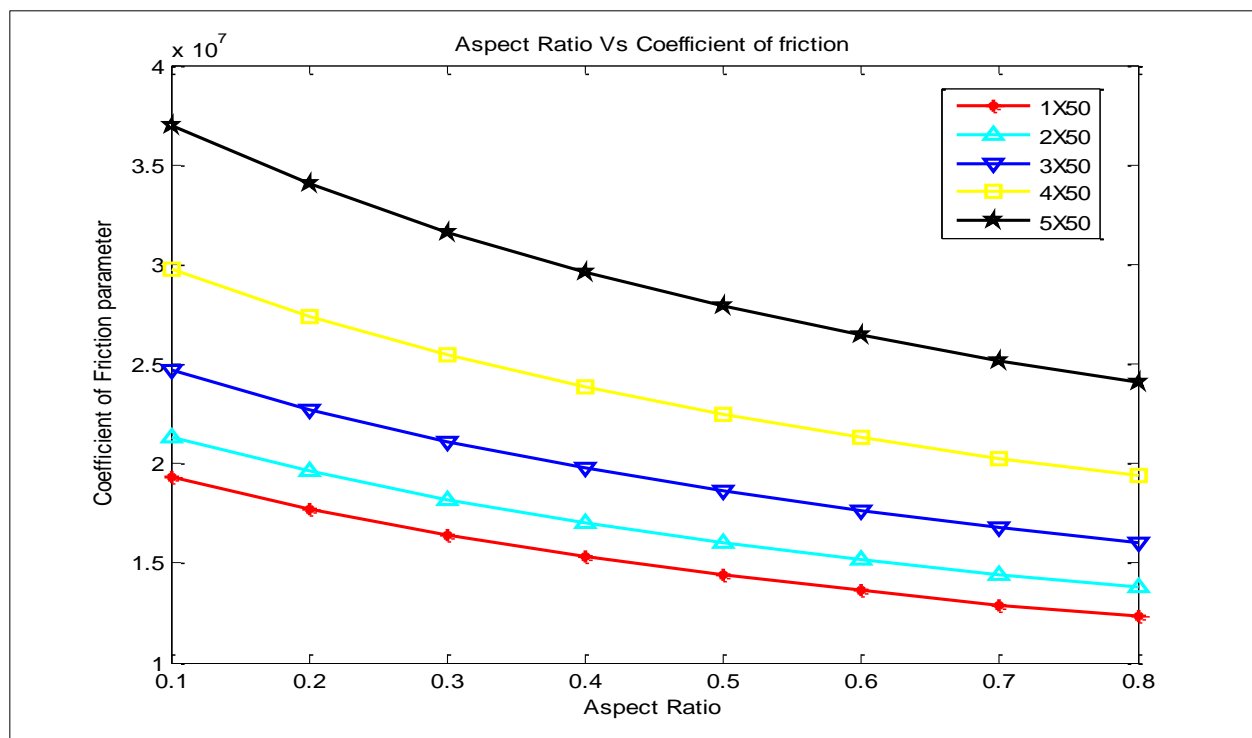


Fig 4.19 Aspect ratio Vs Coefficient of friction parameter

As we can see that from the above figure that for a spherically textured hydrodynamic journal bearing, on increasing the aspect ratio (area of dimple to the area of cell), coefficient of friction parameters decreases, and with increase in number of textures on the surface, coefficient of friction parameter increases for e.g. coefficient of friction parameter of surface having textures 5X50 is more than the surface having textures 1X50, so from above we can conclude that, more the number of textures on the surface, more will be the coefficient of friction parameter. Slope of

the curve between asperity ratio and coefficient of friction parameter is negative because coefficient of friction is decreasing as we increase the asperity ratio.

4.13 Variation of non-dimensional load with increase in texture height for different number of textures

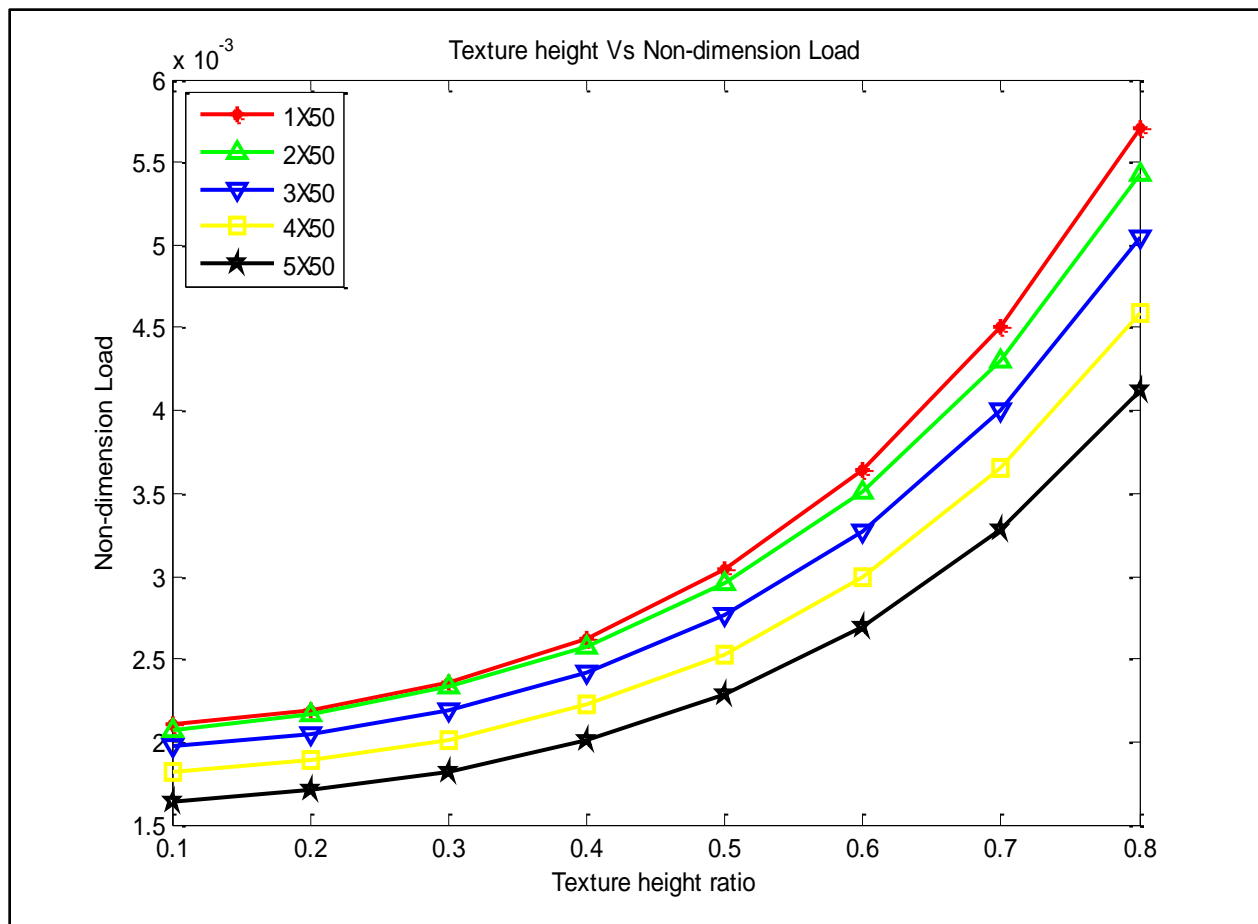


Fig 4.20 Texture height ratio Vs Non dimensional

As we can see from above figure that for a spherically textured journal bearing, on increasing the texture height, non-dimensional load carrying capacity increases and also we noticed from the figure that the surface having high number of textures have slightly low load carrying capacity

than remaining others for e.g. surface having texture 3X50(blue line) have larger load carrying capacity than surface having texture 5X50(black line). Slope of the curve is positive because as we are increase the texture height, non-dimensional load carrying capacity also increases.

4.14 Variation of coefficient of friction parameters with increase in texture height for different number of textures

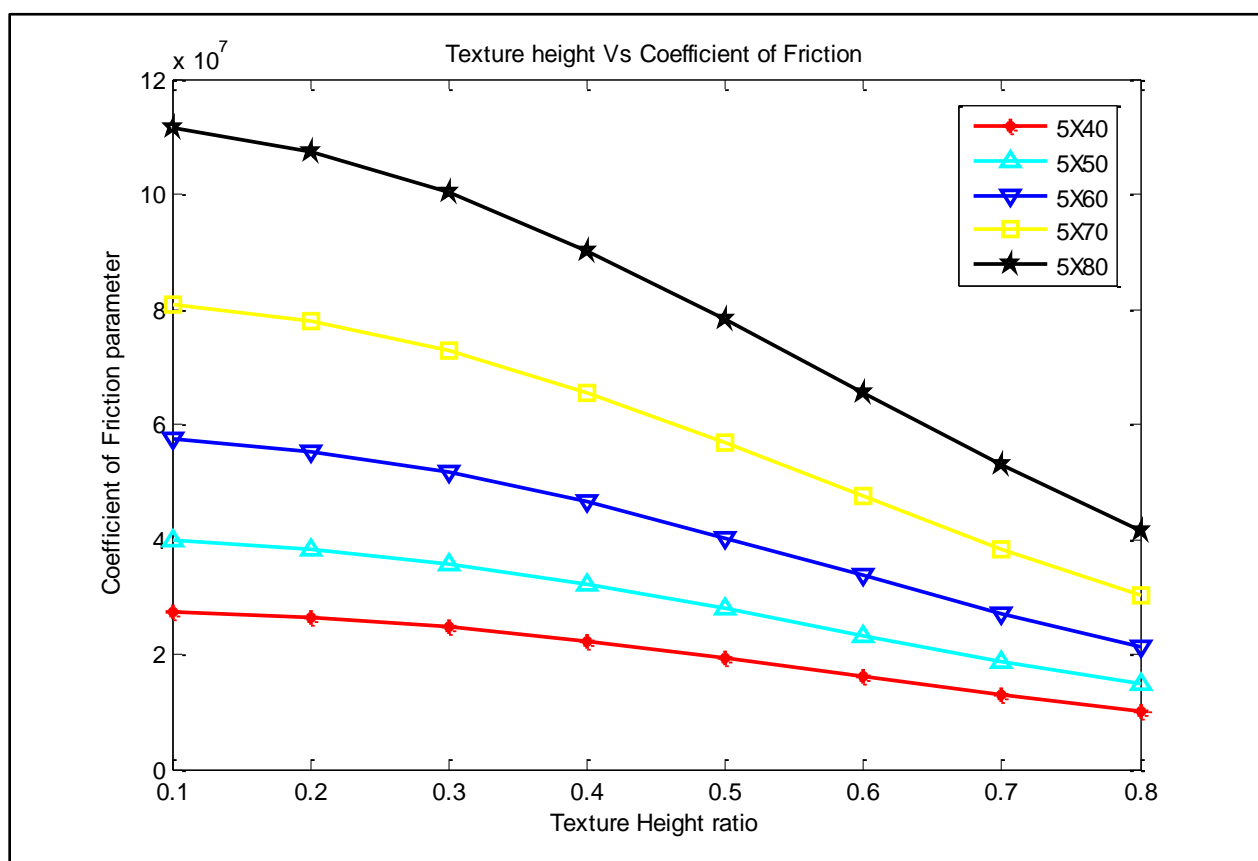


Fig 4.21 Texture height ratio vs Coefficient of friction parameter

From above figure, we can see that for spherically textured hydrodynamic journal bearing, on increasing texture height coefficient of friction parameter decreases because surface appears to be smooth as texture height increases and we also observe from above figure that surface having

higher number of textures on it, having more coefficient of friction parameters than the surface having low number of textures. This is because, surface having higher number of textures, have higher roughness due to the texture portion, which will increase friction for e.g. surface having 5X80(black curve) number of textures have high coefficient of friction parameter than the surface having 5X40(red curve) as shown in figure 4.21. Slope of all curves is negative because as we are increase texture height coefficient of friction parameter decreases.

Conclusion and Future Scope

5.1 Conclusion

- On increasing eccentricity ratio, the non-dimensional load carrying capacity of hydrodynamic journal bearing increases for both spherically textured as well as for normal bearing, but non-dimensional load carrying capacity is more for the former one.
- With increase in eccentricity ratio, the coefficient of friction parameter of hydrodynamic journal bearing decreases for both spherically textured and normal bearing.
- On increasing eccentricity ratio, the non-dimensional flow rate increases for both spherically textured hydrodynamic journal bearing and normal journal bearing but it increases more for the latter one.
- With increase in eccentricity ratio, stiffness of spherically textured journal bearing is slightly less than that of normal journal bearing.
- On increasing eccentricity ratio, damping of spherically textured journal bearing is slightly less than that of normal journal bearing.
- On increasing in asperity ratio, non-dimensional load carrying capacity increases for spherically textured hydrodynamic journal bearing.
- With increase in number of textures, non-dimensional load carrying capacity decreases.
- On increasing asperity ratio, coefficient of friction parameter for spherically textured hydrodynamic journal bearing decreases.
- With increase in number of textures, coefficient of friction parameter increases for spherically textured hydrodynamic journal bearing.

- On increasing texture height, non-dimensional load carrying capacity increases for spherically textured hydrodynamic journal bearing.
- With increase in texture height, coefficient of friction parameter decreases for spherically textured hydrodynamic journal bearing.

5.2 Future scope

- Study the behavior of dynamic characteristics of spherically textured journal bearings by distributed texture partially on the bearing surface.
- To provide spherical texture on piston cylinder arrangement between piston and piston rings and analyze its effect on the dynamic characteristics of the system.
- Study the behavior of dynamic characteristics of journal bearing, by applying different shapes of texture.
- Experimental validation of results.

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